

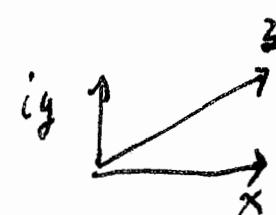
7.6 at max disp (constant), $K = 0$

$$\text{so } E_{\text{tot}} = K + u = U_0$$

and E_{tot} is conserved, so when there is no displacement, $u=0$ and so $K = E_{\text{tot}} = U_0$.

7.11

a.)



$$z = x + iy$$

$$|z|^2 = x^2 + y^2$$

$$zz^* = (x+iy)(x-iy)$$

$$= x^2 + iyx - ixy - i^2 y^2$$

$$= x^2 + y^2 \approx |z|^2$$

b.) $|3w|^2 = (3w)(3w)^* = 3w z^* w^* = 3z^* w w^*$
 $= |z|^2 |w|^2$

thus $|3w| = |z||w|$ — have to take positive square root because $|3w|$ is positive

c.) if $\Phi(x,t) = \psi(x) e^{-i\omega t}$

$$\text{then } |\Phi(x,t)| = |\psi(x)| |e^{-i\omega t}| = |\psi(x)|$$

$$\text{because } |e^{-i\omega t}|^2 = (e^{-i\omega t})(e^{+i\omega t}) = 1$$

7.13 a.) $A \sin(\omega t + \phi) = A(\sin \omega t \cos \phi + \sin \phi \cos \omega t)$
 $= b \sin \omega t + c \cos \omega t$

with $b = A \cos \phi$ and $c = A \sin \phi$

b.) ~~if $\omega t' = \omega t + \phi$~~ the Vice versa:
 $a^2 + b^2 = A^2 (\sin^2 \phi + \cos^2 \phi) = A^2 \Rightarrow A = \pm \sqrt{a^2 + b^2}$

and $\frac{\phi}{b} = \frac{A \sin \phi}{A \cos \phi} = \tan \phi \Rightarrow \phi = \tan^{-1}(a/b)$

(2)

7.26

$$\psi'' = -k^2 \psi : \frac{d^2}{dx^2} \sin kx = \frac{d}{dx} (\cos kx) k = -k^2 \sin kx$$

a.) and $\frac{d^2}{dx^2} \cos kx = -k^2 \cos kx$

and $\frac{d^2}{dx^2} e^{ikx} = -k^2 e^{ikx}$

so they are all solutions, or is e^{-ikx}

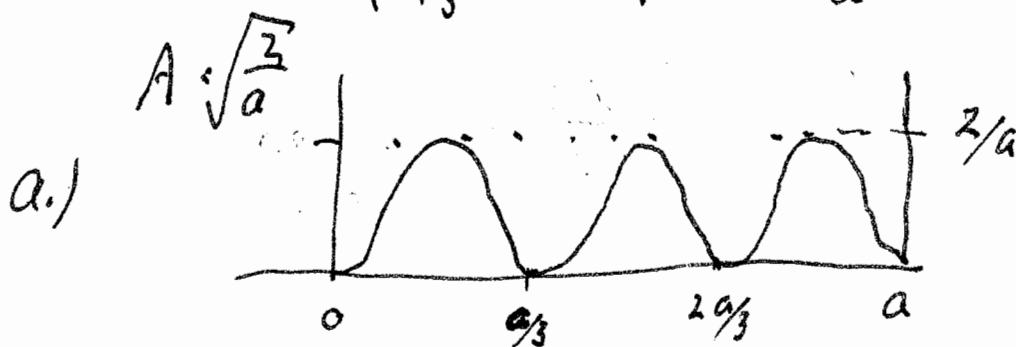
b.) $\sin kx = -i e^{ikx} + a \cos kx$
 $= -i \cos kx + a \cos kx - i^2 \sin kx$
 $= \sin kx \text{ if } a = +i$

so $\sin kx = i [\cos kx \cdot e^{ikx}]$

$\cos kx = e^{ikx} - i \sin kx$

and $e^{ikx} = \cos kx + i \sin kx$

7.30 for $n=3$ $|\psi_3(x)|^2 = |A \sin \frac{3\pi x}{a}|^2$ for box from 0 to a



b.) most probable are where $\sin \frac{3\pi x}{a} = \pm 1$

so $3\pi x/a = \pi/2, 3\pi/2, 5\pi/2$

$x = a\pi/6, a\pi/2, 5a\pi/6$

(3)

both these intervals are
0.01a long. Let's do this
approximately first, then exactly

$$\int_A^{A+\delta} |\psi(x)|^2 dx \doteq |\psi(A+\frac{\delta}{2})|^2 \delta$$

for $A = 0.50a$, $\delta = 0.01a$

$$|\psi(x)|^2 = \frac{2}{a} \sin^2\left(\frac{3\pi x}{a}\right) \Rightarrow \frac{2}{a} \sin^2\left(0.505a \cdot \frac{3\pi}{a}\right)(0.01a)$$

$$= (0.02)(0.998) = 0.0200 = \text{Probability}$$

and for $A = 0.75a$, $\delta = 0.01a$

we get $(0.02)\sin^2(0.755 \cdot 3\pi)$
 $= (0.02)(0.547) = 0.011 = \text{probability}$

to do it exactly we need

$$\int_A^B \frac{2}{a} \sin^2(kx) dx \quad \text{with } k = 3\pi/a$$

$$\text{Note } \cos 2\theta = \cos^2 \theta - \sin^2 \theta$$

$$\text{and } \cos^2 \theta + \sin^2 \theta = 1$$

$$\text{so } \cos 2\theta = 1 - 2 \sin^2 \theta$$

$$\text{or } \sin^2 \theta = (1 - \cos 2\theta)/2$$

So the integral is $\frac{1}{a} \int_A^B [1 - \cos(2kx)] dx$

$$= \frac{1}{a}(B-A) - \frac{1}{a} \frac{1}{2k} (\sin(2kB) - \sin(2kA))$$

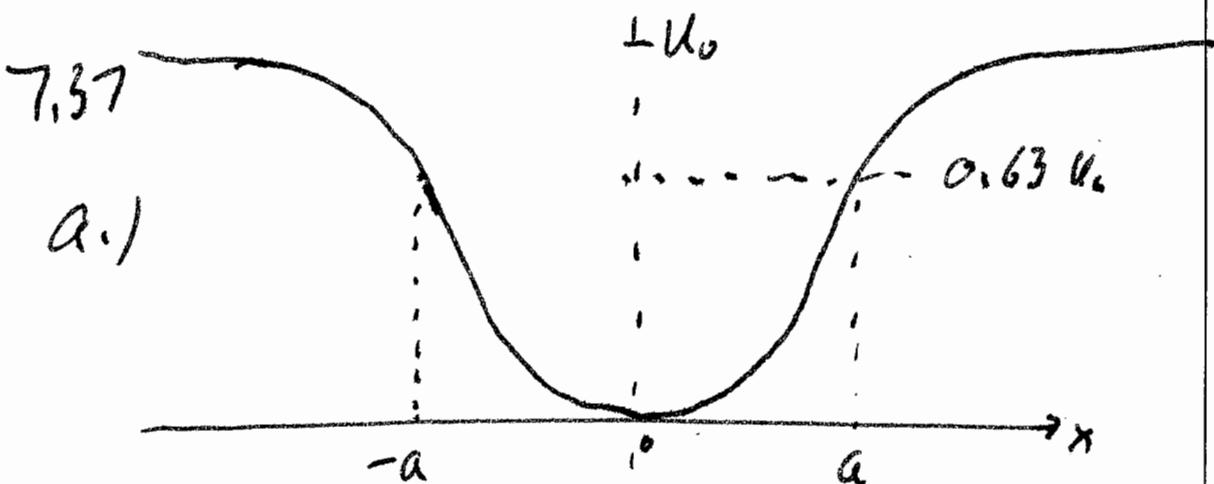
$$= \frac{1}{a}(0.01a) - \frac{1}{6\pi} (\sin 6\pi \left[\frac{51}{76}\right] - \sin 6\pi \left[\frac{50}{75}\right])$$

$$\approx 0.01 + \frac{1}{6\pi} \left([0] - \left[\begin{smallmatrix} -0.187 \\ 0.981 \end{smallmatrix} \right] \right)$$

$$= 0.01 + 0.0099 = 0.0199 \text{ for first interval}$$

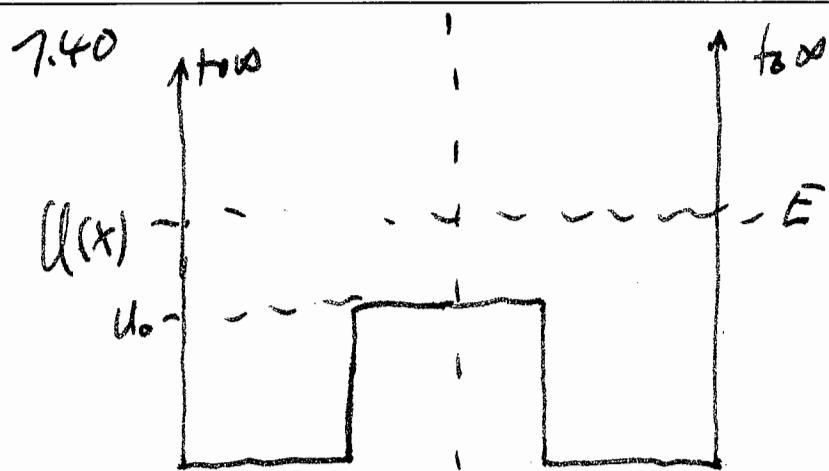
$$- 0.0010 = 0.0110 \text{ for second}$$

The exact result for this small interval is nearly the same as the approximate one.

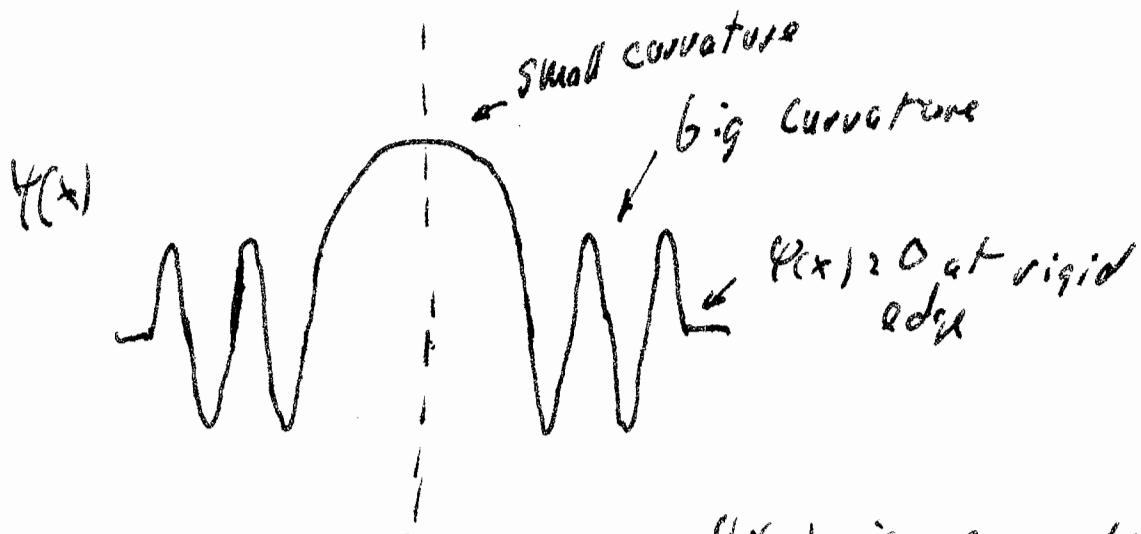


- b.) Classical turning happens where
- $$E = U(x) = U_0 - U_0 e^{-x^2/a^2}$$
- so $U_0 - E = U_0 e^{-x^2/a^2}$
- Take ln: $\ln(U_0 - E) = \ln U_0 + \ln e^{-x^2/a^2}$
- $$\ln U_0 - \ln(U_0 - E) = +x^2/a^2 \quad \text{at turning}$$
- $$x = \pm a \left(\frac{\ln U_0}{\ln(U_0 - E)} \right)^{1/2} \quad \text{at turning}$$

- 7.40 this wave function should be symmetric about the middle of the well, and the k value is bigger (λ shorter) at the edges where $U(x) = 0$. Whether or not any nodes are in the region where $U(x) = U_0$ depends on E, U_0, etc



$E > k$, where
 E would be
 depends on mass
 and well dimension.



Line
of Symmetry

$u(x)$ is symmetric
 about this line
 if number of nodes
 was odd, it would
 be antisymmetric

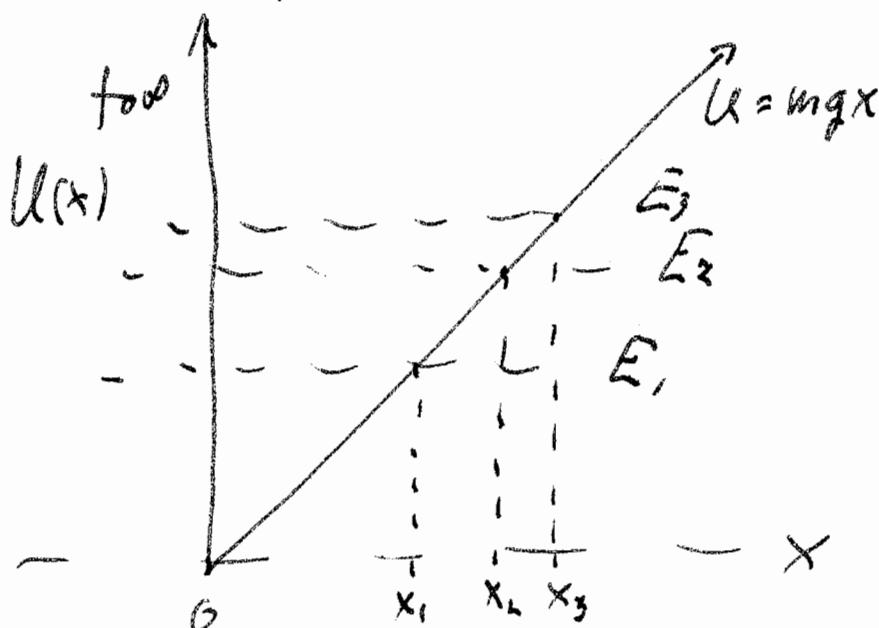
the maximum at
 $x=0$ may be bigger

or ~~smaller~~ smaller

than the others, depending on c_{10}/c_{11} .

for uniform grav. field $U(h) = \text{high}$

and we can call the height x instead of h . For the "hard surface" we can assume it is rigid, and have $U(0) \rightarrow \infty$.



classical turning at $x_i = E_i/mg$



ground state -
no nodes except at
rigid floor

