## AMath 483/583 - Lecture 23

Outline:

- Linear systems: LU factorization and condition number
- Heat equation and discretization
- Iterative methods

Sample codes:

- \$UWHPSC/codes/openmp/jacobi1d_omp1.f90
- \$UWHPSC/codes/openmp/jacobi1d_omp2.f90


## Announcements

Homework 6 is in the notes and due next Friday.
Quizzes for this week's lectures due next Wednesday.
Office hours today 9:30-10:20.

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Next week:
Monday: no class
Wednesday: Guest lecture -

Brad Chamberlain, Cray<br>Chapel: A Next-Generation Partitioned Global Address Space (PGAS) Language

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## DGESV - Solves a general linear system

$$
\begin{array}{ll}
\text { SUBROUTINE DGESV ( } \mathrm{N}, \mathrm{NRHS}, \mathrm{~A}, ~ L D A, ~ I P I V, \\
\& & \mathrm{~B}, \mathrm{LDB}, \mathrm{INFO})
\end{array}
$$

NRHS = number of right hand sides
$B=$ matrix whose columns are right hand side(s) on input solution vector(s) on output.

LDB $=$ leading dimension of $B$.
INFO $=$ integer returning 0 if successful.
$A=$ matrix on input, $L, U$ factors on output,
IPIV = Returns pivot vector (permutation of rows)
integer, dimension(N)
Row I was interchanged with row IPIV (I).

## Gaussian elimination as factorization

If $A$ is nonsingular it can be factored as

$$
P A=L U
$$

where
$P$ is a permutation matrix (rows of identity permuted),
$L$ is lower triangular with 1's on diagonal,
$U$ is upper triangular.
After returning from dgesv:
A contains $L$ and $U$ (without the diagonal of $L$ ), IP IV gives ordering of rows in $P$.

## Gaussian elimination as factorization

Example:

$$
A=\left[\begin{array}{lll}
2 & 1 & 3 \\
4 & 3 & 6 \\
2 & 3 & 4
\end{array}\right]
$$

$\left[\begin{array}{lll}0 & 1 & 0 \\ 0 & 0 & 1 \\ 1 & 0 & 0\end{array}\right]\left[\begin{array}{lll}2 & 1 & 3 \\ 4 & 3 & 6 \\ 2 & 3 & 4\end{array}\right]=\left[\begin{array}{ccc}1 & 0 & 0 \\ 1 / 2 & 1 & 0 \\ 1 / 2 & -1 / 3 & 1\end{array}\right]\left[\begin{array}{ccc}4 & 3 & 6 \\ 0 & 1.5 & 1 \\ 0 & 0 & 1 / 3\end{array}\right]$
$\operatorname{IPIV}=(2,3,1)$
and A comes back from DGESV as:

$$
\left[\begin{array}{ccc}
4 & 3 & 6 \\
1 / 2 & 1.5 & 1 \\
1 / 2 & -1 / 3 & 1 / 3
\end{array}\right]
$$

## dgesv examples

See \$UWHP SC/codes/lapack/random.
Sample codes that solve the linear system $A x=b$ with a random $n \times n$ matrix $A$, where the value $n$ is run-time input.
randomsys1.f90 is with static array allocation.
randomsys $2 . f 90$ is with dynamic array allocation.

## dgesv examples

See \$UWHP SC / codes/lapack/random.
Sample codes that solve the linear system $A x=b$ with a random $n \times n$ matrix $A$, where the value $n$ is run-time input.
randomsys1.f90 is with static array allocation.
randomsys $2 . f 90$ is with dynamic array allocation.
randomsys $3 . \mathrm{f} 90$ also estimates condition number of $A$.

$$
\kappa(A)=\|A\|\left\|A^{-1}\right\|
$$

Can bound relative error in solution in terms of relative error in data using this:

$$
A x^{*}=b^{*} \text { and } A \tilde{x}=\tilde{b} \Longrightarrow \frac{\left\|\tilde{x}-x^{*}\right\|}{\left\|x^{*}\right\|} \leq \kappa(A) \frac{\left\|\tilde{b}-b^{*}\right\|}{\left\|b^{*}\right\|}
$$

## Heat Equation / Diffusion Equation

Partial differential equation (PDE) for $u(x, t)$
in one space dimension and time.
$u$ represents temperature in a 1-dimensional metal rod.
Or concentration of a chemical diffusing in a tube of water.

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The PDE is

$$
u_{t}(x, t)=D u_{x x}(x, t)+f(x, t)
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where subscripts represent partial derivatives,
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$D=$ diffusion coefficient (assumed constant in space \& time),
$f(x, t)=$ source term (heat or chemical being added/removed).
Also need initial conditions $u(x, 0)$ and boundary conditions $u\left(x_{1}, t\right), u\left(x_{2}, t\right)$.

## Steady state diffusion

If $f(x, t)=f(x)$ does not depend on time and if the boundary conditions don't depend on time, then $u(x, t)$ will converge towards steady state distribution satisfying

$$
0=D u_{x x}(x)+f(x)
$$

(by setting $u_{t}=0$.)
This is now an ordinary differential equation (ODE) for $u(x)$.

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$$

(by setting $u_{t}=0$.)
This is now an ordinary differential equation (ODE) for $u(x)$.
We can solve this on an interval, say $0 \leq x \leq 1$ with
Boundary conditions:

$$
u(0)=\alpha, \quad u(1)=\beta
$$

## Steady state diffusion

More generally: Take $D=1$ or absorb in $f$,

$$
u_{x x}(x)=-f(x) \quad \text { for } 0 \leq x \leq 1
$$

Boundary conditions:

$$
u(0)=\alpha, \quad u(1)=\beta
$$

Can be solved exactly if we can integrate $f$ twice and use boundary conditions to choose the two constants of integration.

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Example: $\alpha=20, \beta=60, \quad f(x)=0$ (no heat source)
Solution: $u(x)=\alpha+x(\beta-\alpha) \quad \Longrightarrow u^{\prime \prime}(x)=0$.
No heat source $\Longrightarrow$ linear variation in steady state $\left(u_{x x}=0\right)$.

## Steady state diffusion

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$$
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$$

Boundary conditions:

$$
u(0)=\alpha, \quad u(1)=\beta
$$

Can be solved exactly if we can integrate $f$ twice and use boundary conditions to choose the two constants of integration.

More interesting example:
Example: $\alpha=20, \beta=60, \quad f(x)=100 e^{x}$,
Solution: $u(x)=(100 e-60) x+120-100 e^{x}$.

## Steady state diffusion



## Steady state diffusion



For more complicated equations, numerical methods must generally be used, giving approximations at discrete points.

## Finite difference method

Define grid points $x_{i}=i \Delta x$ in interval $0 \leq x \leq 1$, where

$$
\Delta x=\frac{1}{n+1}
$$

So $x_{0}=0, x_{n+1}=1$, and the $n$ grid points $x_{1}, x_{2}, \ldots, x_{n}$ are equally spaced inside the interval.

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Let $U_{i} \approx u\left(x_{i}\right)$ denote approximate solution.
We know $U_{0}=\alpha$ and $U_{n+1}=\beta$ from boundary conditions.

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Let $U_{i} \approx u\left(x_{i}\right)$ denote approximate solution.
We know $U_{0}=\alpha$ and $U_{n+1}=\beta$ from boundary conditions.
Idea: Replace differential equation for $u(x)$ by system of $n$ algebraic equations for $U_{i}$ values $(i=1,2, \ldots, n)$.

## Finite difference method

$$
\begin{aligned}
& U_{i} \approx u\left(x_{i}\right) \\
& u_{x}\left(x_{i+1 / 2}\right) \approx \frac{U_{i+1}-U_{i}}{\Delta x} \\
& u_{x}\left(x_{i-1 / 2}\right) \approx \frac{U_{i}-U_{i-1}}{\Delta x}
\end{aligned}
$$

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\end{aligned}
$$

So we can approximate second derivative at $x_{i}$ by:

$$
\begin{aligned}
u_{x x}\left(x_{i}\right) & \approx \frac{1}{\Delta x}\left(\frac{U_{i+1}-U_{i}}{\Delta x}-\frac{U_{i}-U_{i-1}}{\Delta x}\right) \\
& =\frac{1}{\Delta x^{2}}\left(U_{i-1}-2 U_{i}+U_{i+1}\right)
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\end{aligned}
$$

This gives coupled system of $n$ linear equations:

$$
\frac{1}{\Delta x^{2}}\left(U_{i-1}-2 U_{i}+U_{i+1}\right)=-f\left(x_{i}\right)
$$

for $i=1,2, \ldots, n$. With $U_{0}=\alpha$ and $U_{n+1}=\beta$.

## Tridiagonal linear system

$$
\begin{aligned}
\alpha-2 U_{1}+U_{2} & =-\Delta x^{2} f\left(x_{1}\right) & (i=1) \\
U_{1}-2 U_{2}+U_{3} & =-\Delta x^{2} f\left(x_{2}\right) & (i=2)
\end{aligned}
$$

Etc.
For $n=5$ :

$$
\left[\begin{array}{ccccc}
-2 & 1 & 0 & 0 & 0 \\
1 & -2 & 1 & 0 & 0 \\
0 & 1 & -2 & 1 & 0 \\
0 & 0 & 1 & -2 & 1 \\
0 & 0 & 0 & 1 & -2
\end{array}\right]\left[\begin{array}{c}
U_{1} \\
U_{2} \\
U_{3} \\
U_{4} \\
U_{5}
\end{array}\right]=-\Delta x^{2}\left[\begin{array}{c}
f\left(x_{1}\right) \\
f\left(x_{2}\right) \\
f\left(x_{3}\right) \\
f\left(x_{4}\right) \\
f\left(x_{5}\right)
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\alpha \\
0 \\
0 \\
0 \\
\beta
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\end{array}\right]-\left[\begin{array}{c}
\alpha \\
0 \\
0 \\
0 \\
\beta
\end{array}\right] .
$$

General $n \times n$ system requires $O\left(n^{3}\right)$ flops to solve.
Tridiagonal $n \times n$ system requires $O(n)$ flops to solve.
Could use LAPACK routine dgtsv.

## Heat equation in 2 dimensions

One-dimensional equation generalizes to

$$
u_{t}(x, y, t)=D\left(u_{x x}(x, y, t)+u_{y y}(x, y, t)\right)+f(x, y, t)
$$

on some domain in the $x-y$ plane, with initial and boundary conditions.

We will only consider rectangle $0 \leq x \leq 1, \quad 0 \leq y \leq 1$.

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Steady state problem (with $D=1$ ):

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This is a PDE in two spatial variables. (Poisson Problem)

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This is a PDE in two spatial variables. (Poisson Problem)
Laplace's equation if $f(x, y) \equiv 0$.
$\nabla^{2}=\left(\partial_{x}^{2}+\partial_{y}^{2}\right)$ is the Laplacian operator.

## Finite difference equations for 2D Poisson problem

## Let $U_{i j} \approx u\left(x_{i}, y_{j}\right)$.

Replace differential equation

$$
u_{x x}(x, y)+u_{y y}(x, y)=-f(x, y)
$$

by algebraic equations

$$
\begin{aligned}
& \frac{1}{\Delta x^{2}}\left(U_{i-1, j}-2 U_{i, j}+U_{i+1, j}\right) \\
& \quad \quad+\frac{1}{\Delta y^{2}}\left(U_{i, j-1}-2 U_{i, j}+U_{i, j+1}\right)=-f\left(x_{i}, y_{j}\right)
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\end{aligned}
$$

If $\Delta x=\Delta y=h$ :

$$
\frac{1}{h^{2}}\left(U_{i-1, j}+U_{i+1, j}+U_{i, j-1}+U_{i, j+1}-4 U_{i, j}\right)=-f\left(x_{i}, y_{j}\right)
$$

## Finite difference equations for 2D Poisson problem

$$
\frac{1}{h^{2}}\left(U_{i-1, j}+U_{i+1, j}+U_{i, j-1}+U_{i, j+1}-4 U_{i, j}\right)=-f\left(x_{i}, y_{j}\right)
$$

On $n \times n \operatorname{grid}(\Delta x=\Delta y=1 /(n+1))$ this gives a linear system of $n^{2}$ equations in $n^{2}$ unknowns.

The above equation must be satisfied for $i=1,2, \ldots, n$ and $j=1,2, \ldots, n$.

Matrix is $n^{2} \times n^{2}$,
e.g. on 100 by 100 grid, matrix is $10,000 \times 10,000$.

Contains $(10,000)^{2}=100,000,000$ elements.

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Matrix is $n^{2} \times n^{2}$,
e.g. on 100 by 100 grid, matrix is $10,000 \times 10,000$.

Contains $(10,000)^{2}=100,000,000$ elements.
Matrix is sparse: each row has at most 5 nonzeros out of $n^{2}$ elements! But structure is no longer tridiagonal.

## Finite difference equations for 2D Poisson problem



Matrix has block tridiagonal structure:

$$
A=\frac{1}{h^{2}}\left[\begin{array}{cccc}
T & I & & \\
I & T & I & \\
& I & T & I \\
& & I & T
\end{array}\right] \quad T=\left[\begin{array}{rrrr}
-4 & 1 & & \\
1 & -4 & 1 & \\
& 1 & -4 & 1 \\
& & 1 & -4
\end{array}\right]
$$

## Iterative methods

Back to one space dimension first...
Coupled system of $n$ linear equations:

$$
\left(U_{i-1}-2 U_{i}+U_{i+1}\right)=-\Delta x^{2} f\left(x_{i}\right)
$$

for $i=1,2, \ldots, n$. With $U_{0}=\alpha$ and $U_{n+1}=\beta$.
Iterative method starts with initial guess $U^{[0]}$ to solution and then improves $U^{[k]}$ to get $U^{[k+1]}$ for $k=0,1, \ldots$.

Note: Generally does not involve modifying matrix $A$.
Do not have to store matrix $A$ at all, only know about stencil.

## Jacobi iteration

$$
\left(U_{i-1}-2 U_{i}+U_{i+1}\right)=-\Delta x^{2} f\left(x_{i}\right)
$$

Solve for $U_{i}$ :

$$
U_{i}=\frac{1}{2}\left(U_{i-1}+U_{i+1}+\Delta x^{2} f\left(x_{i}\right)\right)
$$

Note: With no heat source, $f(x)=0$, the temperature at each point is average of neighbors.

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Note: With no heat source, $f(x)=0$,
the temperature at each point is average of neighbors.
Suppose $U^{[k]}$ is a approximation to solution. Set

$$
U_{i}^{[k+1]}=\frac{1}{2}\left(U_{i-1}^{[k]}+U_{i+1}^{[k]}+\Delta x^{2} f\left(x_{i}\right)\right) \text { for } i=1,2, \ldots, n .
$$

Repeat for $k=0,1,2, \ldots$ until convergence.
Can be shown to converge (eventually... very slow!)

## Slow convergence of Jacobi



## Slow convergence of Jacobi



## Slow convergence of Jacobi



## Iterative methods

Jacobi iteration is about the worst possible iterative method.
But it's very simple, and useful as a test for parallelization.

Better iterative methods:

- Gauss-Seidel
- Successive Over-Relaxation (SOR)
- Conjugate gradients
- Preconditioned conjugate gradients
- Multigrid


## Iterative methods - initialization

```
! allocate storage for boundary points too:
allocate (x \((0: n+1), u(0: n+1), f(0: n+1))\)
\(d x=1 \cdot d 0 /(n+1 . d 0)\)
! \$omp parallel do
do \(i=0, n+1\)
    ! grid points:
    \(x(i)=i * d x\)
    ! source term:
    f(i) = 100.*exp(x(i))
    ! initial guess (linear function):
    \(u(i)=a l p h a+x(i) *(b e t a-a l p h a)\)
    enddo
```


## Jacobi iteration in Fortran

```
uold = u ! starting values before updating
do iter=1,maxiter
    dumax = 0.d0
    do i=1,n
        u(i) = 0.5d0*(uold(i-1) + uold(i+1) + dx**2*f(i))
        dumax = max(dumax, abs(u(i)-uold(i)))
        enddo
    ! check for convergence:
    if (dumax .lt. tol) exit
    uold = u ! for next iteration
    enddo
```

Note: we must use old value at $i-1$ for Jacobi.
Otherwise we get the Gauss-Seidel method.

$$
u(i)=0.5 d 0 *(u(i-1)+u(i+1)+d x * * 2 * f(i))
$$

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```

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Otherwise we get the Gauss-Seidel method.

$$
u(i)=0.5 d 0 *(u(i-1)+u(i+1)+d x * * 2 * f(i))
$$

This actually converges faster!

## Jacobi with OpenMP parallel do (fine grain)

See: \$UWHPSC/codes/openmp/jacobi1d_omp1.f90

```
uold = u ! starting values before updating
do iter=1,maxiter
dumax = 0.d0
!$omp parallel do reduction(max : dumax)
do i=1,n}= = 0.5d0*(uold(i-1) + uold(i+1) + dx**2*f(i)
    dumax = max(dumax, abs(u(i)-uold(i)))
    enddo
! check for convergence:
if (dumax.lt. tol) exit
!$omp parallel do
do i=1,n
    uold(i) = u(i) ! for next iteration
enddo
```

Note: Forking threads twice each iteration.

## Jacobi with OpenMP - coarse grain

General Approach:

- Fork threads only once at start of program.
- Each thread is responsible for some portion of the arrays, from $i=i s t a r t$ to $i=i e n d$.
- Each iteration, must copy u to uold, update u, check for convergence.
- Convergence check requires coordination between threads to get global dumax.
- Print out final result after leaving parallel block

See code in the repository or the notes:
\$UWHPSC/codes/openmp/jacobi1d_omp2.f90

