### Conservation Laws and Finite Volume Methods AMath 574 Winter Quarter, 2017

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http://faculty.washington.edu/rjl/classes/am574w2017

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- Riemann solution for shallow water: http://faculty.washington.edu/rjl/classes/am574w2017/riemann.html
- Euler equations of gas dynamics
- Notebooks from https://github.com/clawpack/riemann\_book

Reading: Chapter 14

## Compressible gas dynamics

In one space dimension (e.g. in a pipe).  $\rho(x,t) = \text{density}, \quad u(x,t) = \text{velocity},$  $p(x,t) = \text{pressure}, \quad \rho(x,t)u(x,t) = \text{momentum}.$ 

Conservation of:

Conservation laws:

$$\rho_t + (\rho u)_x = 0$$
$$(\rho u)_t + (\rho u^2 + p)_x = 0$$

Equation of state:

$$p = P(\rho).$$

(Later: p may also depend on internal energy / temperature)

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### Conservation laws:



Momentum flux:

 $\rho u^2 = (\rho u)u = advective flux$ 

p term in flux?

- $-p_x =$  force in Newton's second law,
- as momentum flux: microscopic motion of gas molecules.

## Momentum flux arising from pressure



Note that:

- molecules with positive *x*-velocity crossing  $x_1$  to right increase the momentum in  $[x_1, x_2]$
- molecules with negative *x*-velocity crossing *x*<sub>1</sub> to left also increase the momentum in [*x*<sub>1</sub>, *x*<sub>2</sub>]

Hence momentum flux increases with pressure  $p(x_1, t)$  even if macroscopic (average) velocity is zero.

## Compressible gas dynamics

Conservation laws:

$$\rho_t + (\rho u)_x = 0$$
$$(\rho u)_t + (\rho u^2 + p)_x = 0$$

Equation of state:

$$p = P(\rho).$$

Same as shallow water if  $P(\rho) = \frac{1}{2}g\rho^2$  (with  $\rho \equiv h$ ).

Isothermal:  $P(\rho) = a^2 \rho$  (since *T* proportional to  $p/\rho$ ). Isentropic:  $P(\rho) = \hat{\kappa} \rho^{\gamma}$  ( $\gamma \approx 1.4$  for air)

Jacobian matrix:

$$f'(q) = \begin{bmatrix} 0 & 1\\ P'(\rho) - u^2 & 2u \end{bmatrix}, \qquad \lambda = u \pm \sqrt{P'(\rho)}.$$

Dam break problem for shallow water equations with tracer

$$h_t + (hu)_x = 0$$
$$(hu)_t + \left(hu^2 + \frac{1}{2}gh^2\right)_x = 0$$
$$(h\phi)_t + (uh\phi)_x = 0$$



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# Riemann solution for the SW equations in x-t plane



#### Similarity solution:

Solution is constant on any ray: q(x,t) = Q(x/t)

Riemann solution can be calculated for many problems. Linear: Eigenvector decomposition. Nonlinear: more difficult.

In practice "approximate Riemann solvers" used numerically.

## Euler equations of gas dynamics

Conservation of mass, momentum, energy:  $q_t + f(q)_x = 0$  with

$$q = \begin{bmatrix} \rho \\ \rho u \\ E \end{bmatrix}, \qquad f(q) = \begin{bmatrix} \rho u \\ \rho u^2 + p \\ u(E+p) \end{bmatrix}$$

where  $E = \rho e + \frac{1}{2}\rho u^2$ 

Equation of state:  $p = \text{pressure} = p(\rho, E)$ 

Ideal gas, polytropic EOS:  $p = \rho e(\gamma - 1) = (\gamma - 1) \left(E - \frac{1}{2}\rho u^2\right)$ 

 $\gamma\approx7/5=1.4$  for air,  $\quad\gamma=5/3$  for monatomic gas

The Jacobian f'(q) has eigenvalues u - c, u, u + c where

$$c = \sqrt{rac{dp}{d
ho}} \bigg|_{
m at \ constant \ entropy} = \sqrt{rac{\gamma p}{
ho}} \ {
m for \ ideal \ gas}$$

Initial data:

$$q(x,0) = \begin{cases} q_l & \text{if } x < 0\\ q_r & \text{if } x > 0 \end{cases}$$

Shock tube problem:  $u_l = u_r = 0$ , jump in  $\rho$  and p.

$$\left( \begin{array}{ccc} \rho_{i} & p_{i} & u_{i} \end{array} \right) \left( \begin{array}{ccc} \rho_{r} & \rho_{r} & u_{r} \end{array} \right)$$

Pressure:



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Pressure:



## **Riemann Problem for gas dynamics**

Waves propagating in x-t space:



Similarity solution (function of x/t alone).

#### Waves can be approximated by discontinuties: Approximate Riemann solvers