Conservation Laws and Finite Volume Methods AMath 574 Winter Quarter, 2017

Randall J. LeVeque Applied Mathematics University of Washington

http://faculty.washington.edu/rjl/classes/am574w2017

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- Linear acoustics
- Diagonalization of linear systems
- Meaning of eigenvectors
- Characteristic solution for acoustics
- Riemann problem for acoustics

Reading: Chapter 3

Linear acoustics

Example: Linear acoustics in a 1d gas tube

$$q = \left[\begin{array}{c} p \\ u \end{array} \right] \qquad \begin{array}{c} p(x,t) = \text{pressure perturbation} \\ u(x,t) = \text{velocity} \end{array}$$

Equations:

 $p_t + \kappa u_x = 0$ Change in pressure due to compression $\rho u_t + p_x = 0$ Newton's second law, F = ma

where K = bulk modulus, and $\rho =$ unperturbed density of gas. Hyperbolic system:

$$\left[\begin{array}{c}p\\u\end{array}\right]_t + \left[\begin{array}{cc}0&\kappa\\1/\rho&0\end{array}\right] \left[\begin{array}{c}p\\u\end{array}\right]_x = 0.$$

Linear acoustics

$$\left[\begin{array}{c}p\\u\end{array}\right]_t+\left[\begin{array}{c}0&\kappa\\1/\rho&0\end{array}\right]\left[\begin{array}{c}p\\u\end{array}\right]_x=0.$$

This has the form $q_t + Aq_x = 0$ with

eigenvalues: $\lambda^1 = -c, \qquad \lambda^2 = +c,$

where $c = \sqrt{\kappa/\rho} =$ speed of sound.

eigenvectors:
$$r^1 = \begin{bmatrix} -Z \\ 1 \end{bmatrix}$$
, $r^2 = \begin{bmatrix} Z \\ 1 \end{bmatrix}$

where $Z = \rho c = \sqrt{\rho \kappa} = \text{impedance}.$

$$R = \begin{bmatrix} -Z & Z \\ 1 & 1 \end{bmatrix}, \qquad R^{-1} = \frac{1}{2Z} \begin{bmatrix} -1 & Z \\ 1 & Z \end{bmatrix}$$

.

Special initial data:

$$q(x,0) = \begin{cases} q_l & \text{if } x < 0\\ q_r & \text{if } x > 0 \end{cases}$$

Example: Acoustics with bursting diaphram ($u_l = u_r = 0$)



Pressure:



Acoustic waves propagate with speeds $\pm c$.

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Riemann Problem for acoustics

Waves propagating in x-t space:



Left-going wave $W^1 = q_m - q_l$ and right-going wave $W^2 = q_r - q_m$ are eigenvectors of A.

Consider constant coefficient linear system $q_t + Aq_x = 0$.

Suppose hyperbolic:

- Real eigenvalues $\lambda^1 \leq \lambda^2 \leq \cdots \leq \lambda^m$,
- Linearly independent eigenvectors r^1, r^2, \ldots, r^m .

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Let $R = [r^1 | r^2 | \cdots | r^m]$ $m \times m$ matrix of eigenvectors.

Then $Ar^p = \lambda^p r^p$ means that $AR = R\Lambda$ where

$$\Lambda = \begin{bmatrix} \lambda^1 & & & \\ & \lambda^2 & & \\ & & \ddots & \\ & & & \lambda^m \end{bmatrix} \equiv \operatorname{diag}(\lambda^1, \lambda^2, \dots, \lambda^m).$$

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 $AR = R\Lambda \implies A = R\Lambda R^{-1}$ and $R^{-1}AR = \Lambda$. Similarity transformation with *R* diagonalizes *A*.

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Consider constant coefficient linear system $q_t + Aq_x = 0$. Multiply system by R^{-1} :

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Use $R^{-1}AR = \Lambda$ and define $w(x,t) = R^{-1}q(x,t)$:

 $w_t(x,t) + \Lambda w_x(x,t) = 0.$ Since *R* is constant!

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This decouples to m independent scalar advection equations:

$$w_t^p(x,t) + \lambda^p w_x^p(x,t) = 0.$$
 $p = 1, 2, ..., m.$

Suppose
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 for $-\infty < x < \infty$.

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The solution to the decoupled equation $w_t^p + \lambda^p w_x^p = 0$ is

$$w^{p}(x,t) = w^{p}(x - \lambda^{p}t, 0) = \overset{\circ p}{w}(x - \lambda^{p}t).$$

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We can rewrite this as

$$q(x,t) = \sum_{p=1}^{m} w^{p}(x,t) r^{p} = \sum_{p=1}^{m} w^{op}(x-\lambda^{p}t) r^{p}$$

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Riemann Problem for acoustics



Then

$$q_m = \boldsymbol{w_r^1 r^1} + \boldsymbol{w_l^2 r^2}$$

So the waves W^1 and W^2 are eigenvectors of *A*:

$$\mathcal{W}^1 = q_m - q_l = (w_r^1 - w_l^1)r^1$$

 $\mathcal{W}^2 = q_r - q_m = (w_r^2 - w_l^2)r^2$.

$$q(x,0) = \begin{bmatrix} \overset{\circ}{p}(x) \\ 0 \end{bmatrix} = -\frac{\overset{\circ}{p}(x)}{2Z_0} \begin{bmatrix} -Z_0 \\ 1 \end{bmatrix} + \frac{\overset{\circ}{p}(x)}{2Z_0} \begin{bmatrix} Z_0 \\ 1 \end{bmatrix}$$
$$= w^1(x,0)r^1 + w^2(x,0)r^2$$
$$= \begin{bmatrix} \overset{\circ}{p}(x)/2 \\ -\overset{\circ}{p}(x)/(2Z_0) \end{bmatrix} + \begin{bmatrix} \overset{\circ}{p}(x)/2 \\ \overset{\circ}{p}(x)/(2Z_0) \end{bmatrix}.$$

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Physical meaning of eigenvectors

Eigenvectors for acoustics:

$$r^{1} = \left[\begin{array}{c} -\rho_{0}c_{0} \\ 1 \end{array} \right] = \left[\begin{array}{c} -Z_{0} \\ 1 \end{array} \right], \qquad r^{2} = \left[\begin{array}{c} \rho_{0}c_{0} \\ 1 \end{array} \right] = \left[\begin{array}{c} Z_{0} \\ 1 \end{array} \right]$$

In a simple 1-wave (propagating at speed $\lambda^1 = -c_0$),

$$\left[\begin{array}{c} p_x \\ u_x \end{array}\right] = \beta(x) \left[\begin{array}{c} -Z_0 \\ 1 \end{array}\right]$$

The pressure variation is $-Z_0$ times the velocity variation.

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Similarly, in a simple 2-wave ($\lambda^2 = c_0$),

$$\left[\begin{array}{c} p_x \\ u_x \end{array}\right] = \beta(x) \left[\begin{array}{c} Z_0 \\ 1 \end{array}\right]$$

The pressure variation is Z_0 times the velocity variation.

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Riemann solution for a linear system

Linear hyperbolic system: $q_t + Aq_x = 0$ with $A = R\Lambda R^{-1}$. General Riemann problem data $q_l, q_r \in \mathbb{R}^m$.

Decompose jump in q into eigenvectors:

$$q_r - q_l = \sum_{p=1}^m \alpha^p r^p$$

Note: the vector α of eigen-coefficients is

$$\alpha = R^{-1}(q_r - q_l) = R^{-1}q_r - R^{-1}q_l = w_r - w_l.$$

Riemann solution consists of m waves $\mathcal{W}^p \in \mathbb{R}^m$:

$$\mathcal{W}^p = \alpha^p r^p$$
, propagating with speed $s^p = \lambda^p$.