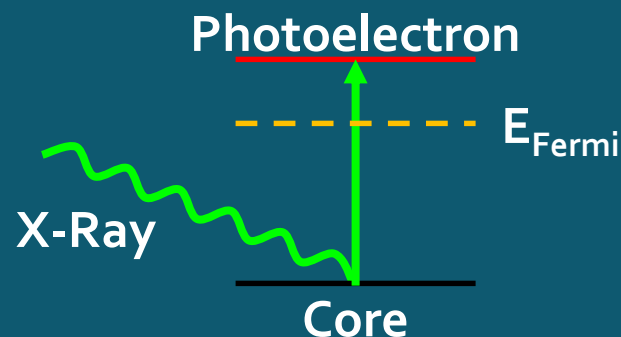
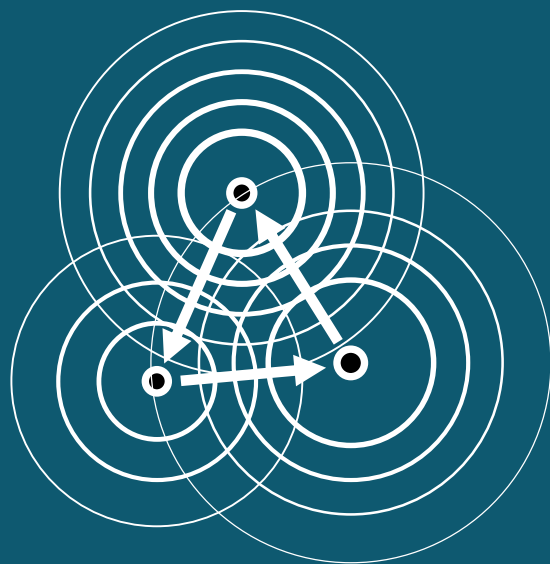


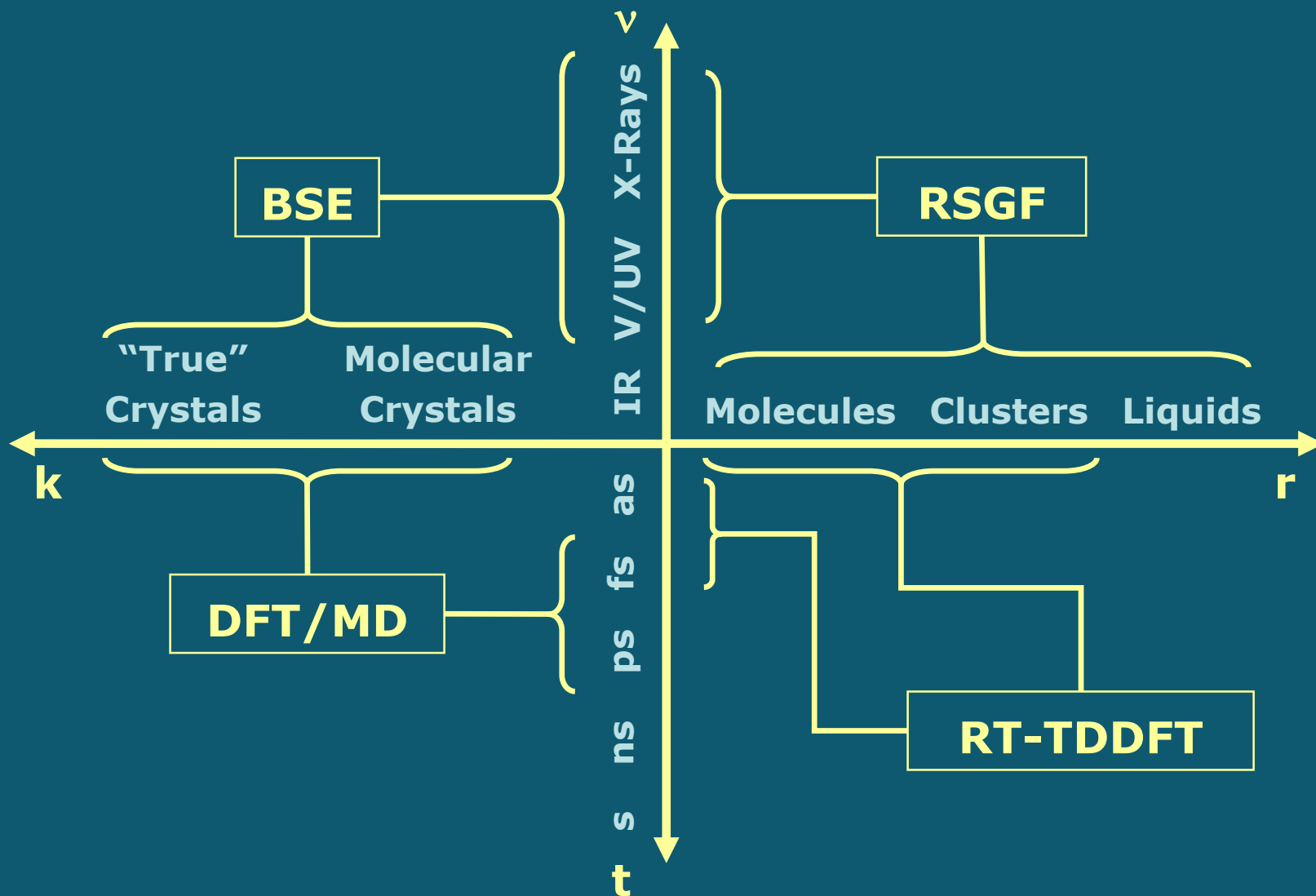
Real-Time and Finite Temperature Green's Function Approaches for Excited States, Response Functions, and Thermodynamics

F. D. Vila, J. J. Kas and J. J. Rehr



DOE Office of Science grants DE-FG02-03ER15476 (FDV) and DE-FG02-97ER45623 (JJK, JJR), with computer support from DOE - NERSC.

Experimental/Theoretical Landscape



Motivation

Why **real-time** and **finite temperature**?

Old approaches **not well suited** for new science:

- XFEL pulsed x-ray sources (FLASH, LCLS)

- Pump-probe experiments (with finite electronic T)

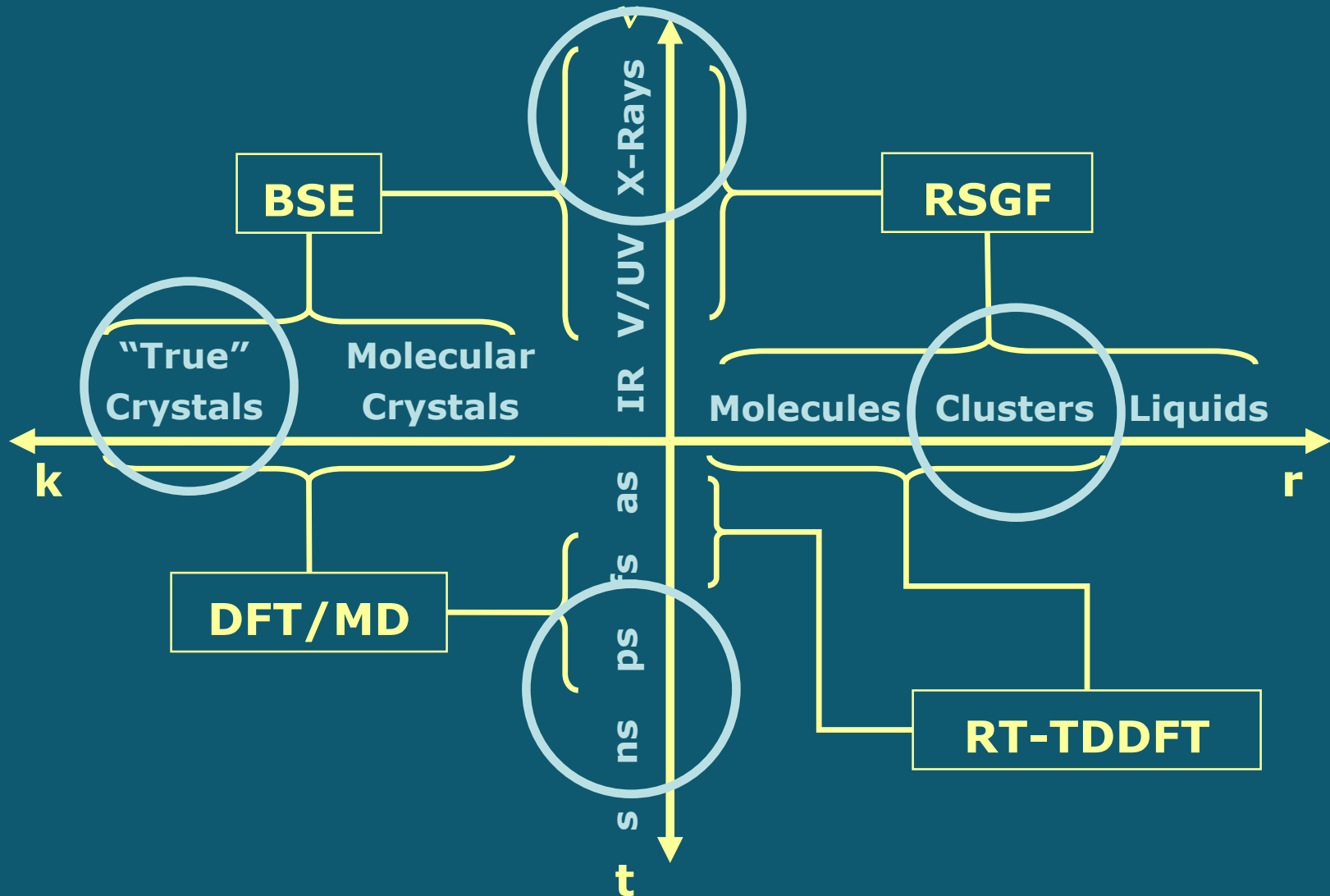
- Interest in **time-dependent** response

- Direct** access to time-domain

- Importance of **non-equilibrium** states

- Interest in **exotic** states (warm dense matter)

The Challenge



This talk:

Green's Functions and DFT/MD Approaches in X-Ray
Excited States and Structural Disorder

Real-Time Approaches for Optical and Core Response

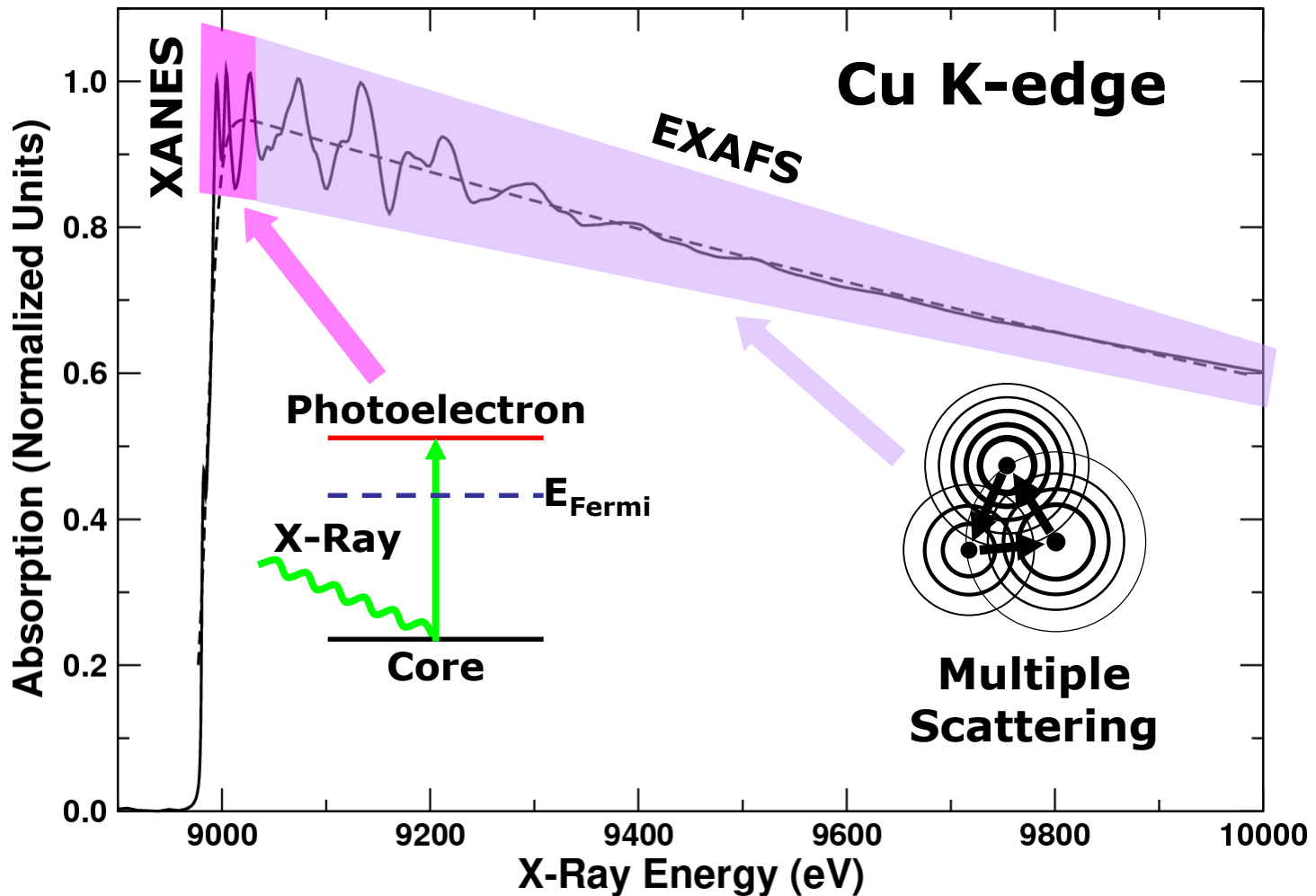
Finite Temperature Green's Functions

Green's Functions and DFT/MD Approaches in X-Ray Excited States and Structural Disorder

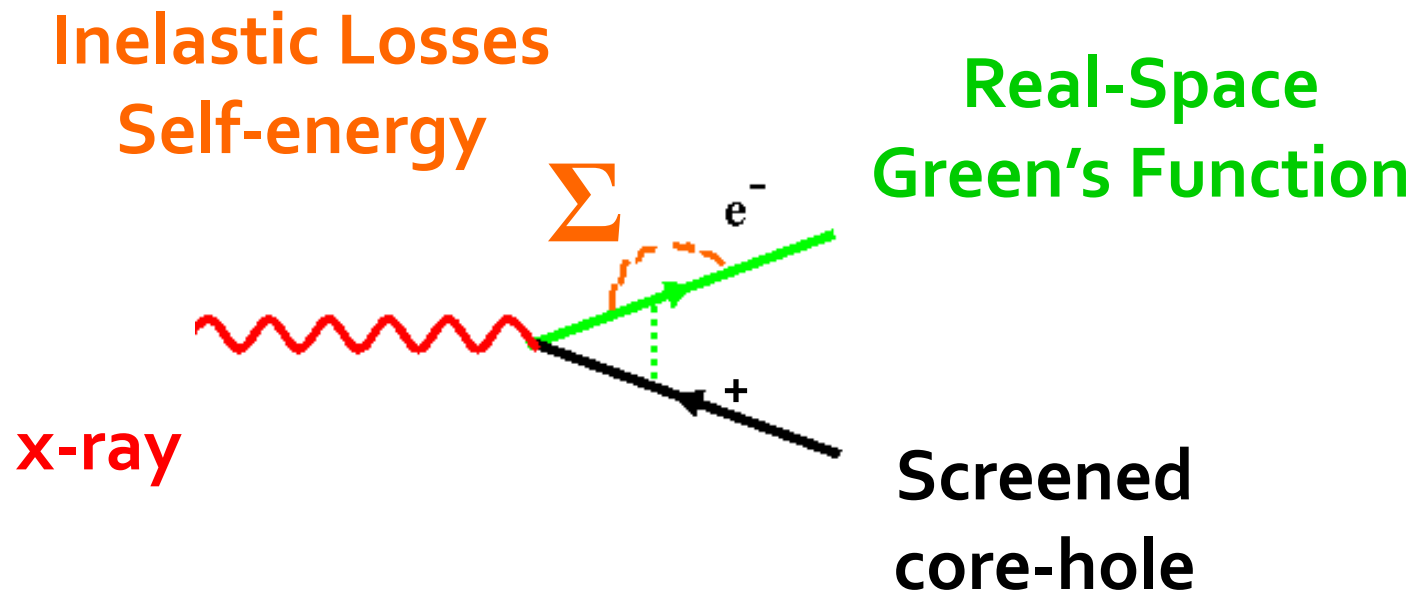
Real-Time Approaches for Optical and Core Response

Finite Temperature Green's Functions

XAS: XANES and EXAFS



FEFF Quantitative XANES Theory in One Picture



Many-body to Effective Single Particle

XAS absorption coefficient

Many-body Fermi's
Golden Rule

$$\mu(\omega) \propto \sum_F |\langle I | \Delta | F \rangle|^2 \delta(E_F - E_I - \omega)$$



Effective Single particle
Fermi's Golden Rule

$$\mu(\omega) \propto \sum_{i,f} |\underbrace{\langle i |}_{\text{H}} \underbrace{d | f \rangle}_{\text{H'}}|^2 \delta(E_f - E_i - \omega)$$

$$H = -(1/2)\nabla^2 + V$$

$$H' = H + V_{ch} + \Sigma(E)$$

FEFF: From Sum-Over-States to Green's Function

**Effective Single particle
Fermi's Golden Rule**

$$\mu(\omega) \propto \sum_{cf} \langle c | d^\dagger | f \rangle \langle f | d | c \rangle \delta(E_f - E_c - \omega)$$

$$\rho(r, r', E) = \sum_f |f\rangle \langle f| \delta(E_f - E)$$

$$\text{Im} [G(r, r', E)] = -\frac{1}{\pi} \rho(r, r', E)$$

$$G = [E - H + i\Gamma]^{-1}$$

**Substitute sum over
final states with
Green's function**

$$\mu(\omega) \propto \text{Im} \sum_c \langle c | d^\dagger G(\omega + E_c) d | c \rangle \theta_\Gamma(\omega + E_c + E_{\text{Fermi}})$$

FEFF: Local basis and Matrix elements

$$\mu(\omega) \propto \text{Im} \sum_c \langle c | d^\dagger G(\omega + E_c) d | c \rangle \theta_\Gamma(\omega + E_c + E_{\text{Fermi}})$$

Insert complete set of **site states** $1 = \sum_L |i, L\rangle \langle i, L|$

Matrix elements

$$\mu(\omega) \propto \text{Im} \sum_{iLL'} \underbrace{\langle c | d^\dagger | i, L \rangle}_{\text{Matrix elements}} \underbrace{G_{LL'}(\omega + E_c)}_{\text{Green's function matrix}} \underbrace{\langle i, L' | d | c \rangle}_{\text{Matrix elements}} \theta_\Gamma(\omega + E_c + E_{\text{Fermi}})$$

Green's function matrix

Getting G: Path Expansion and Full Multiple Scattering

Dyson's equation:

$$G = G^0 + G^0 V G$$

Iterating:

$$G = G^0 + G^0 V G^0 + G^0 V G^0 V G^0 + \dots$$

Atomic pot., site scat.,
central atom, etc.

$$V = \sum_i v_i$$
$$t_i = v_i + v_i G_0 t_i$$

$$G_c = G_0 + G_0 t_c G_0$$

$$T_{LL'}^{ij} = t_{LL'}^i \delta_{ij}$$

EXAFS

Path expansion

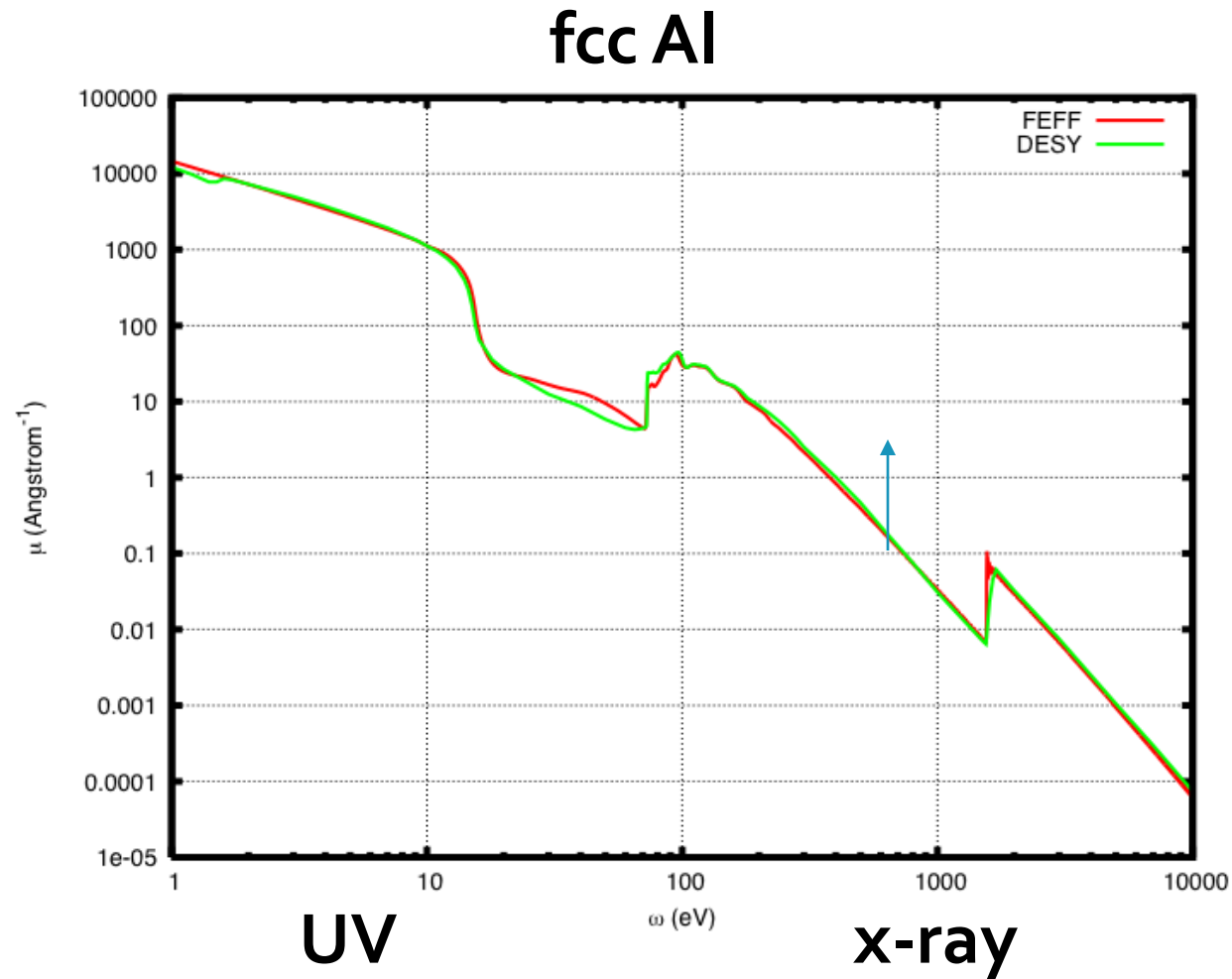

$$G = G_c + G_c T G_c + G_c T G^0 T G_c + G_c T G^0 T G^0 T G_c + \dots$$

XANES

Full Multiple Scattering

$$G = [1 - G^0 T]^{-1} G^0$$

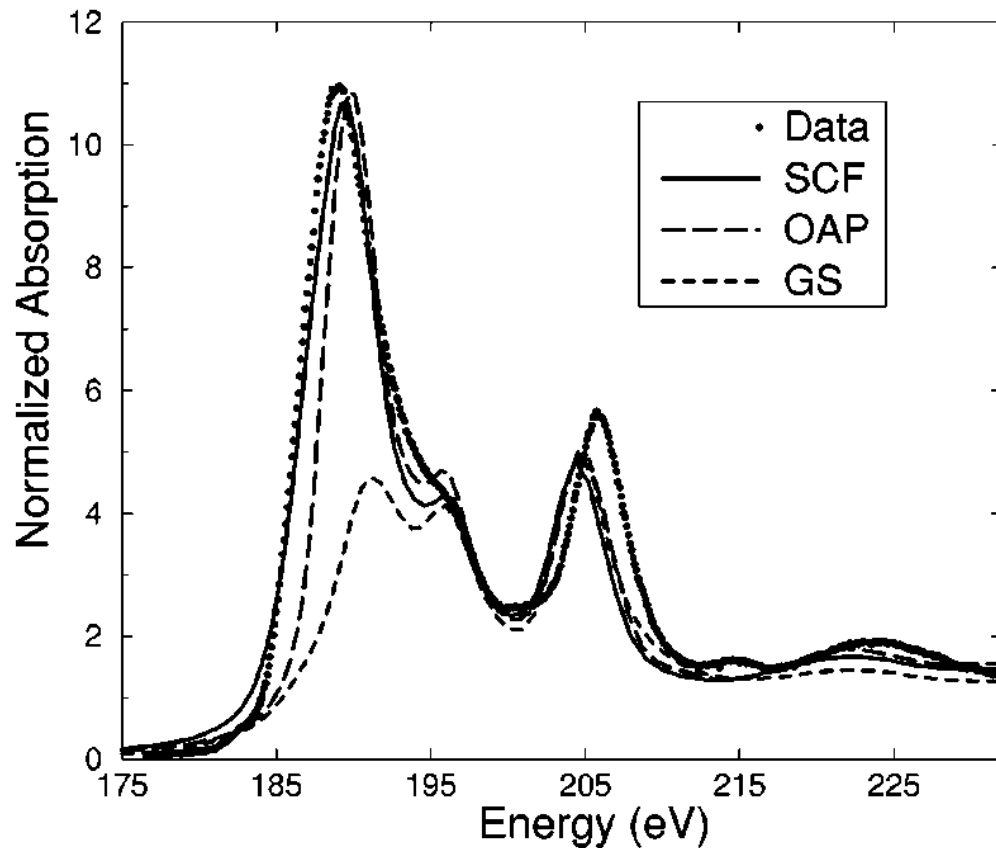
Full spectrum: Expt. Vs Theory



Beyond DFT: Quasi-particle Self-Energy Effects

Quasi-particle (QP) effects:

$$G = [E - H - \Sigma]^{-1}$$



BN 89 atom cluster

Ground state potential:

Usually **insufficient**

Need **QP** effects

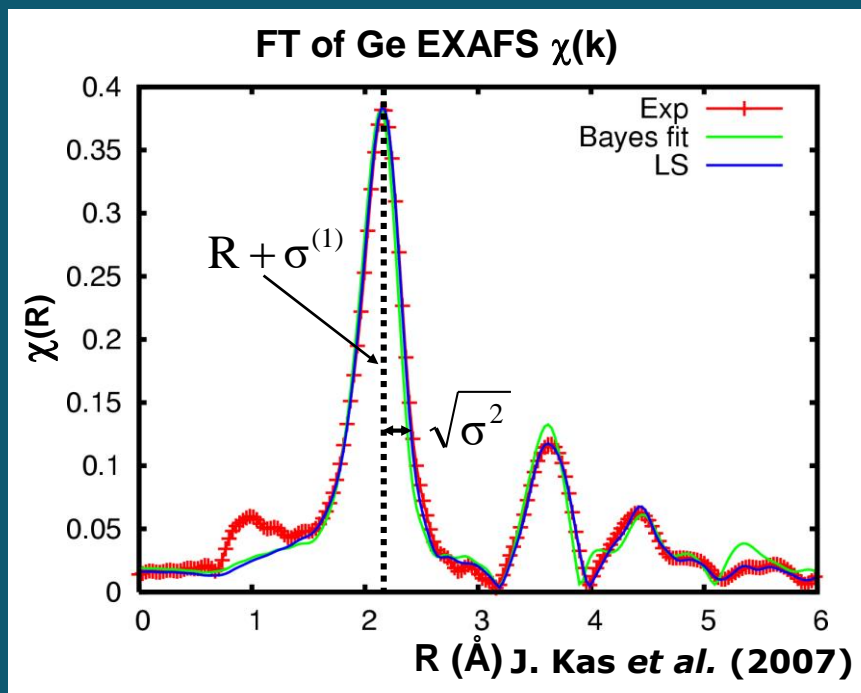
Debye-Waller Factors in EXAFS

Multiple Scattering Path XAFS DW Factor

Average commonly expressed in terms of the **cumulant expansion**

$$\langle e^{i2kr} \rangle \equiv e^{2ikR_0} e^{-W(T)}$$

$$W(T) = - \sum_{n=1}^{\infty} \frac{(2ik)^n}{n!} \sigma^{(n)}(T)$$



First **cumulants**

$$\left\{ \begin{array}{l} \sigma^{(1)} = \langle r - R_0 \rangle \\ \sigma^{(2)} = \langle (r - \bar{r})^2 \rangle \equiv \sigma^2(T) \\ \sigma^{(3)} = \langle (r - \bar{r})^3 \rangle \end{array} \right.$$

DW factors from Phonon Green's Function

XAFS DW Factor for path R:

$$\sigma_R^2(T) = \frac{\hbar}{2\mu_R} \int_0^\infty \frac{1}{\omega} \coth\left(\frac{\beta\hbar\omega}{2}\right) \rho_R(\omega) d\omega$$

VDOS expressed as imaginary part of the **phonon Green's function**

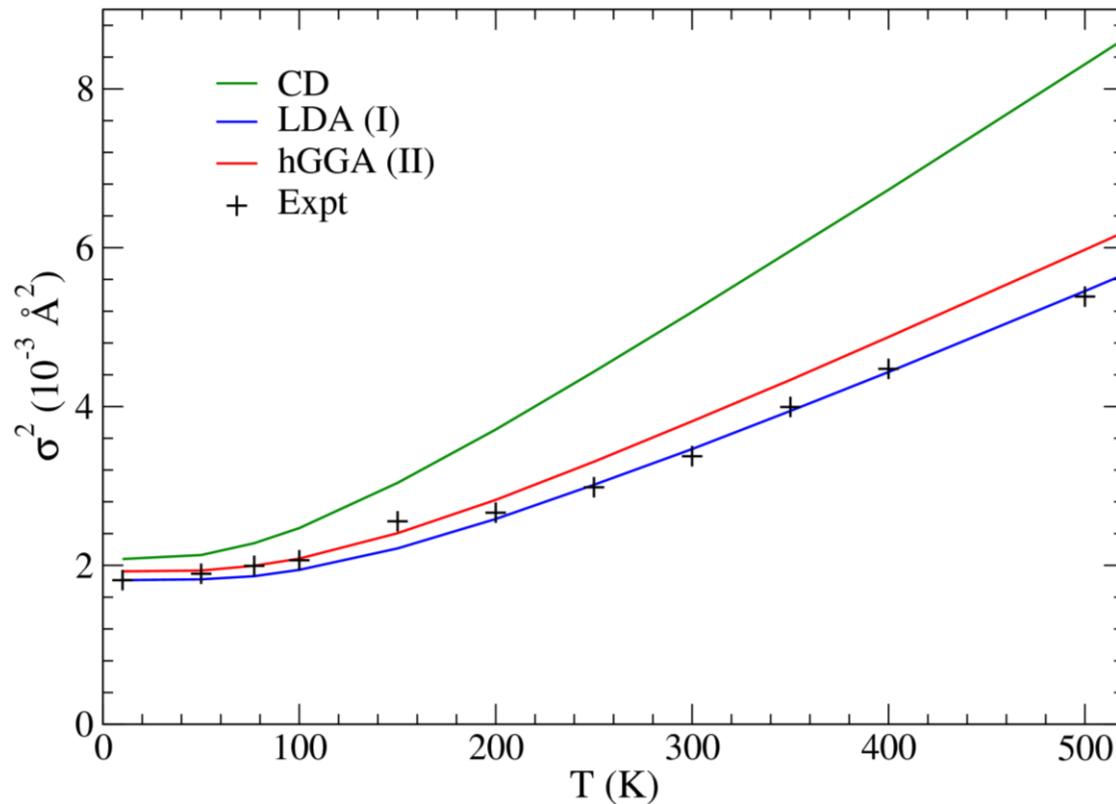
$$\rho_R(\omega) = -\frac{2\omega}{\pi} \text{Im} \left\langle 0 \left| \frac{1}{\omega^2 - \mathbf{D} + i\varepsilon} \right| 0 \right\rangle$$

Seed state: Displacement along path

Dynamical Matrix: Calculated using ab initio methods (Abinit, Gaussian, VASP, NWChem, etc)



EXAFS near-neighbor DW Factor of Ge



CD (Correlated Debye):
Standard FEFF

LDA, hGGA:
Ab initio DW

Directional bonding:
Needs AIDW

Dynamic Structural Disorder In XANES

Dynamic structure in supported Pt nanoclusters: Real-time density functional theory and x-ray spectroscopy simulations

F. D. Vila, J. J. Rehr, J. Kas, R. G. Nuzzo, A. I. Frenkel

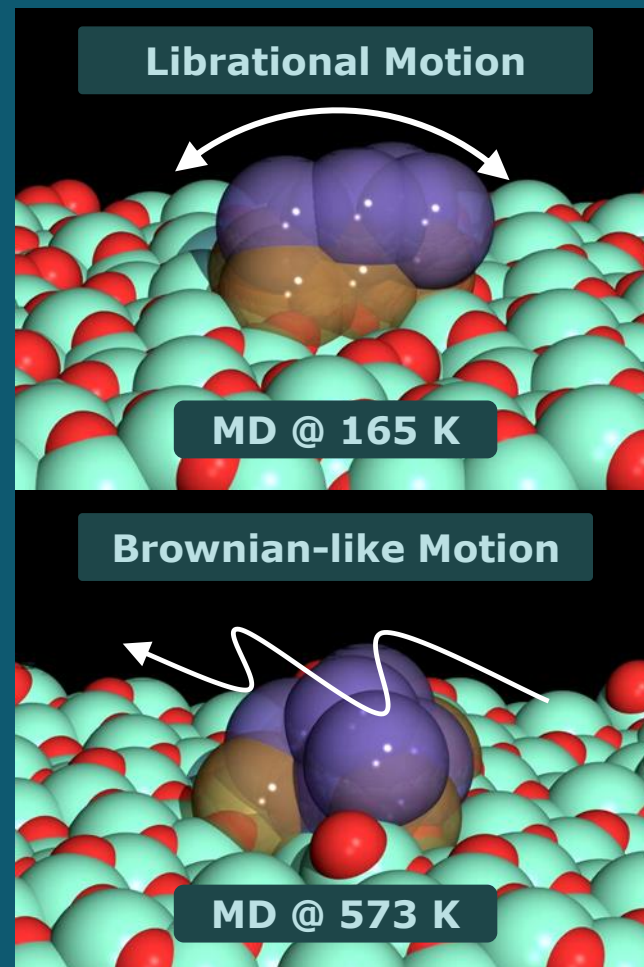
Physical Review B **78**, 121404(R), 2008

Complex dynamics:

multiple-time scales, librational motion, fluctuating bonding

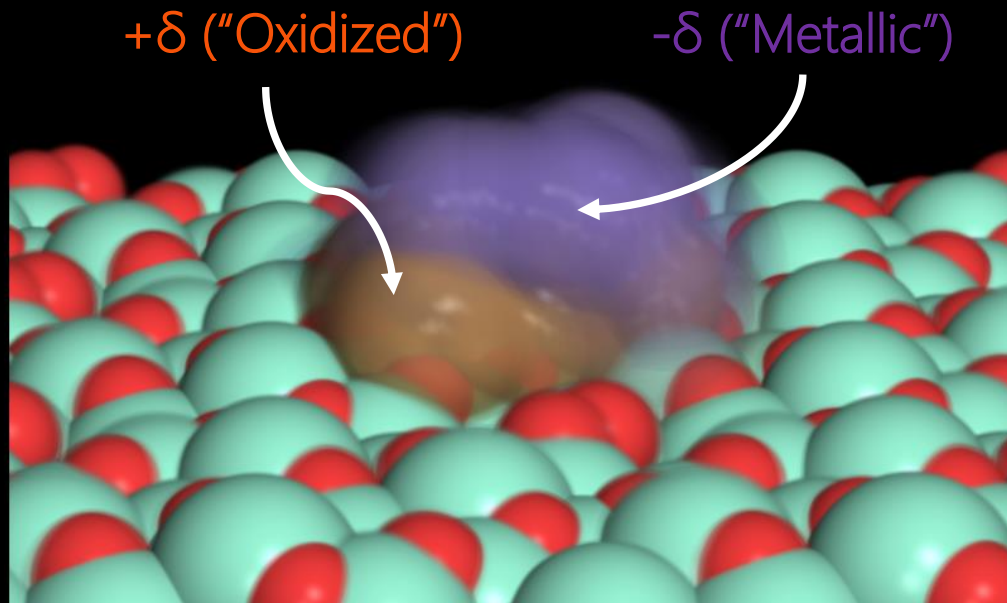
Simulations explain:

large structural **disorder**, Negative Thermal Expansion (**NTE**).

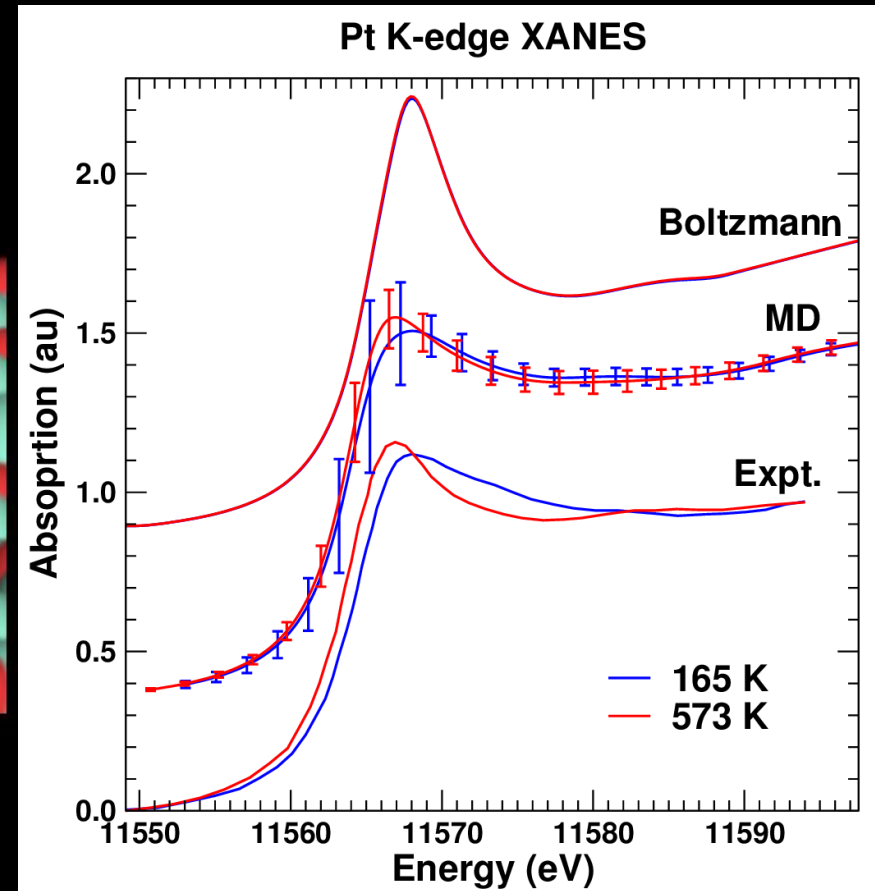


Pt₁₀ on γ -Al₂O₃

Theory: Static Simulations are Inadequate



Pt₁₀ on γ -Al₂O₃ @ 165 K



MD simulations reproduce experiment

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Finite Temperature Green's Functions

Real-Time TDDFT Advantages

Can be **more efficient** than frequency space (large systems)

Very **versatile** (pulses, transport, etc)

More “physical/**realistic**”

Easy access to **NLO properties**

Real-Time Time Dependent DFT

Direct numerical integration of the time-dependent Kohn-Sham equations in a time-dependent external electric field:

$$i\frac{\partial\Psi}{\partial t} = H(t) \Psi \quad \Psi(t) = T \exp\left(-i \int_0^t H(t') dt'\right) \Psi(0)$$
$$H = -\frac{1}{2}\nabla^2 + V_{ext}(\mathbf{r}, t) + V_H[\rho](\mathbf{r}, t) + V_{xc}[\rho](\mathbf{r}, t)$$

Optical properties are determined from the total dipole moment:

$$\mathbf{p}(t) = \int \rho(\mathbf{r}, t) \mathbf{r} d^3\mathbf{r}$$

$$\underbrace{\chi_{ij}^{(1)}(\omega) = \delta p_i(\omega) / E_j(\omega) = \alpha_{ij}(\omega)}_{\text{Linear Response}} \quad \underbrace{\sigma(\omega) \sim \omega \mathbf{Im} \langle \alpha(\omega) \rangle}_{\text{Absorption}}$$

Linear Response

Absorption

Numerical Real-Time Evolution

The **ground state** density ρ_0 , overlap matrix S and Hamiltonian matrix $H(t)$ evaluated at each time-step using **SIESTA**:

$$i\frac{\partial c(t)}{\partial t} = S^{-1}H(t)\underbrace{c(t)}$$

Orbital coefficients

Accurate and **stable** evolution using Crank-Nicolson propagator:

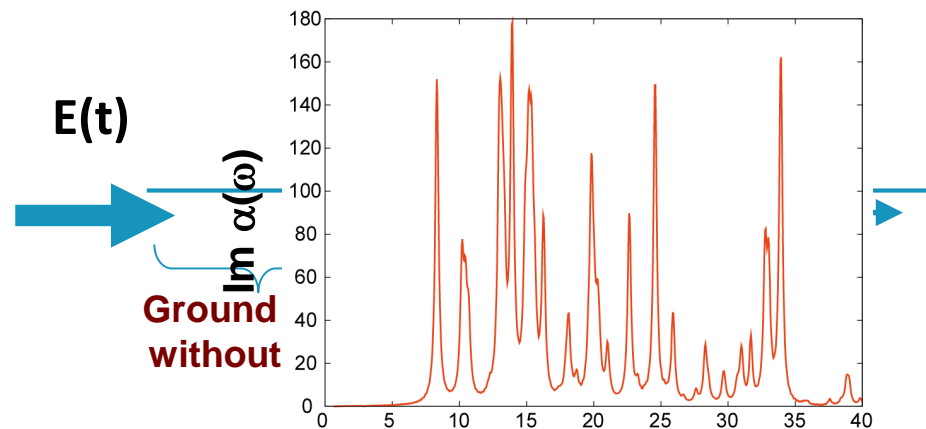
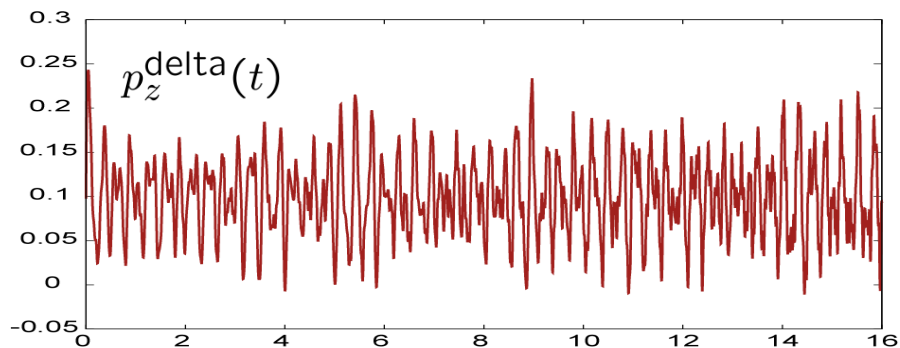
$$c(t + \Delta t) = \frac{1 - iS^{-1}H(\bar{t})\Delta t/2}{1 + iS^{-1}H(\bar{t})\Delta t/2}c(t) + \mathcal{O}(\Delta t^2)$$

Soler et al. J. Phys.: Condens. Matter 14, 2745 (2002)

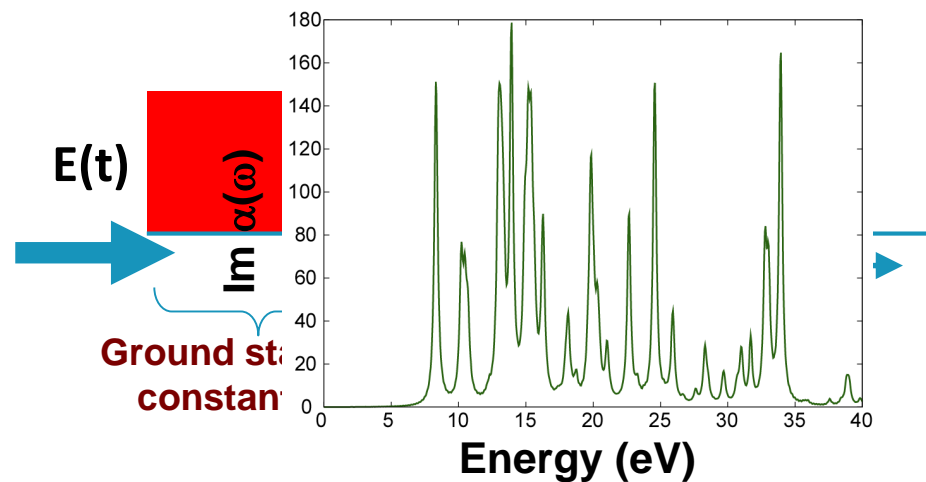
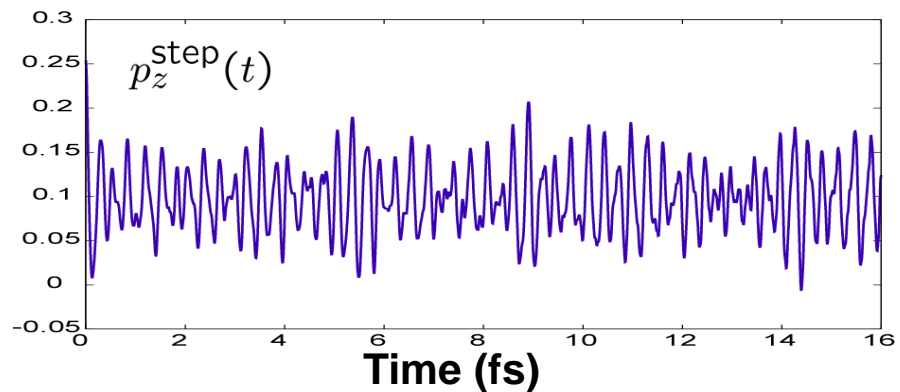
Tsolakidis et al. Phys. Rev. B 66, 235416 (2002)

Optical Absorption in CO

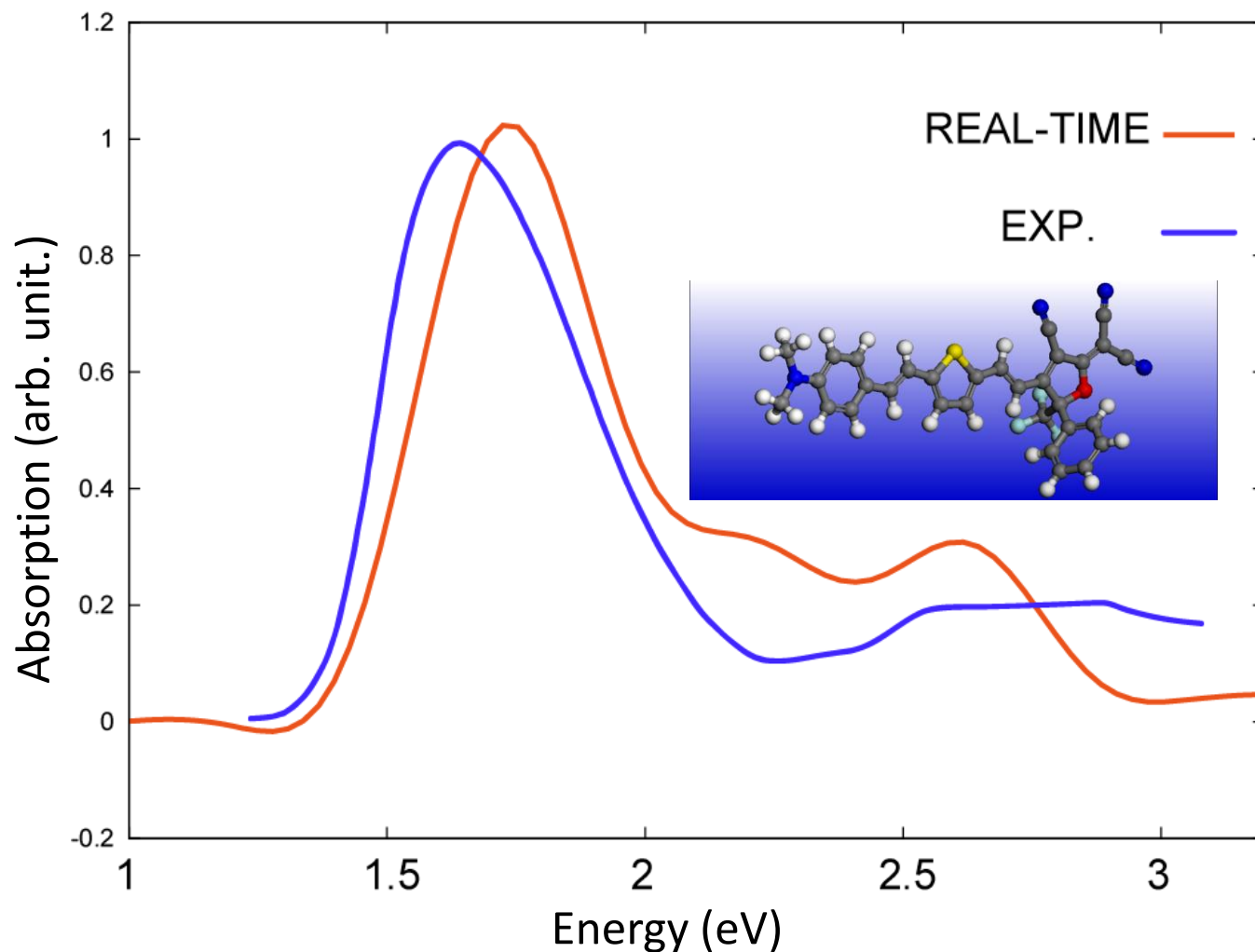
Delta Function (Unit Impulse at $t=0$)



Step Function (Turn-off Constant E at $t=0$)



Linear Response: Chromophore YLD_156



**Real-time (gas)
absorption peak**

$$\omega_0 = 1.72 \text{ eV}$$

$$(\lambda_0 = 721 \text{ nm})$$

**Expt. (CHCl₃)
absorption peak**

$$\omega_0 = 1.65 \text{ eV}$$

$$(\lambda_0 = 753 \text{ nm})$$

Real time Nonlinear Response

Nonlinear expansion:

$$P = \chi^{(1)} E + \chi^{(2)} E^2 + \chi^{(3)} E^3 + \dots$$

Must take into account **response** time:

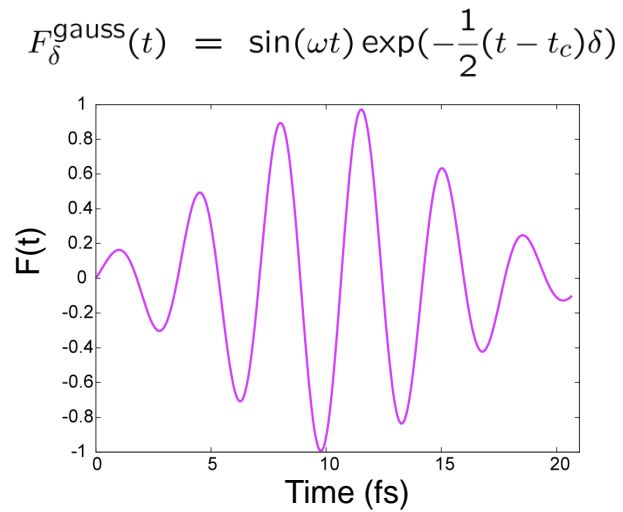
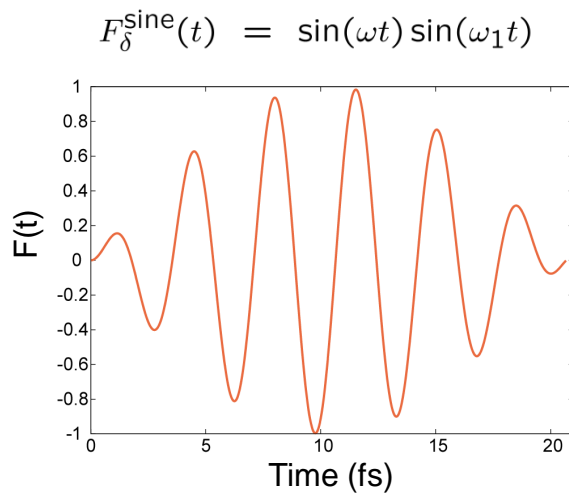
$$\begin{aligned} p_i(t) = & \mu_i^0 + \int dt_1 \chi_{ij}^{(1)}(t - t_1) E_j(t_1) \\ & + \int dt_1 \int dt_2 \chi_{ijk}^{(2)}(t - t_1, t - t_2) E_j(t_1) E_k(t_2) \\ & + \int dt_1 \int dt_2 \int dt_3 \chi_{ijkl}^{(3)}(t - t_1, t - t_2, t - t_3) E_j(t_1) E_k(t_2) E_l(t_3) \\ & + \dots \end{aligned}$$

How do we **invert** to get nonlinear response function?

Dynamic NLO with Quasi-Monochromatic Field F_δ

F_δ : Sine wave with **sine** or **Gaussian envelope**

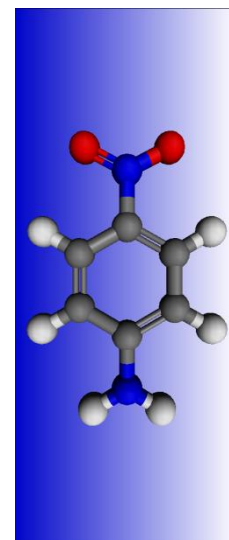
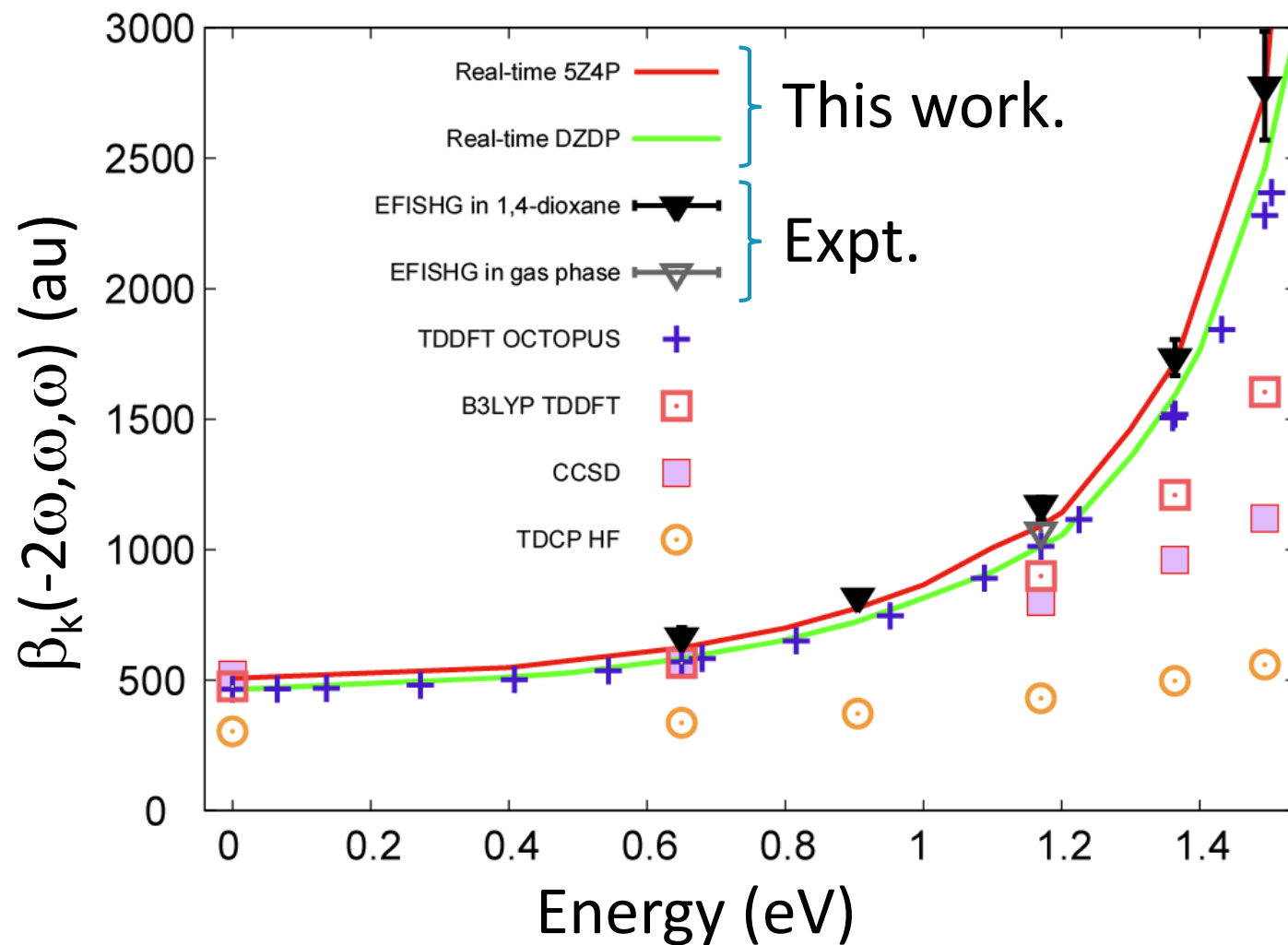
$$\omega = 1.17 \text{ eV}$$
$$\omega_1 = 0.1 \text{ eV}$$



$$\chi_{ijk}^{(2)}(-2\omega_0; \omega_0, \omega_0) = \frac{2\pi p_{ijk}^{(2)}(2\omega_0)}{\int_{-\Delta}^{\Delta} d\omega' F(\omega_0 - \omega') F(\omega_0 + \omega')} \quad \text{SHG}$$

$$\chi_{ijk}^{(2)}(0; -\omega_0, \omega_0) = \frac{\pi p_{ijk}^{(2)}(0)}{\int_{-\Delta}^{\Delta} d\omega' F^*(\omega_0 + \omega') F(\omega_0 + \omega')} \quad \text{OR}$$

Second Harmonic Generation (SHG) in pNA



Real-Time X-Ray Spectroscopy (RTXS)

XAS Absorption (FGR, Δ SCF, FSR)

$$\mu(\omega) = \sum_k |\langle c|d|k\rangle|^2 \delta_\Gamma(\omega + \varepsilon_c - \varepsilon_k) \theta(E - E_F)$$

↓ FT

$$\mu(\omega) = \frac{1}{\pi} \text{Re} \int_0^\infty dt e^{i\omega t} \underline{G_c(t)} \underline{\langle \psi(t) | \psi(0) \rangle} \theta(\omega + \varepsilon_c - E_F)$$

Core Hole Green's Function

$$\underline{G_c(t)} = i \exp[i(\varepsilon_c + i\Gamma)t]$$

Autocorrelation Function

$$\underline{\langle \psi(t) | \psi(0) \rangle} = \sum_{jj'} \langle c|d^\dagger|j\rangle U_{jj'}(t,0) \langle \tilde{j}|d|c\rangle$$

$$\langle \psi(0) | \psi(t) \rangle \begin{cases} |\psi(0)\rangle = d|c\rangle \\ |\psi(t)\rangle = U(t,0)|\psi(0)\rangle \end{cases} \longrightarrow U(t,0) = T \exp \left[-i \int_0^t dt H(t) \right]$$

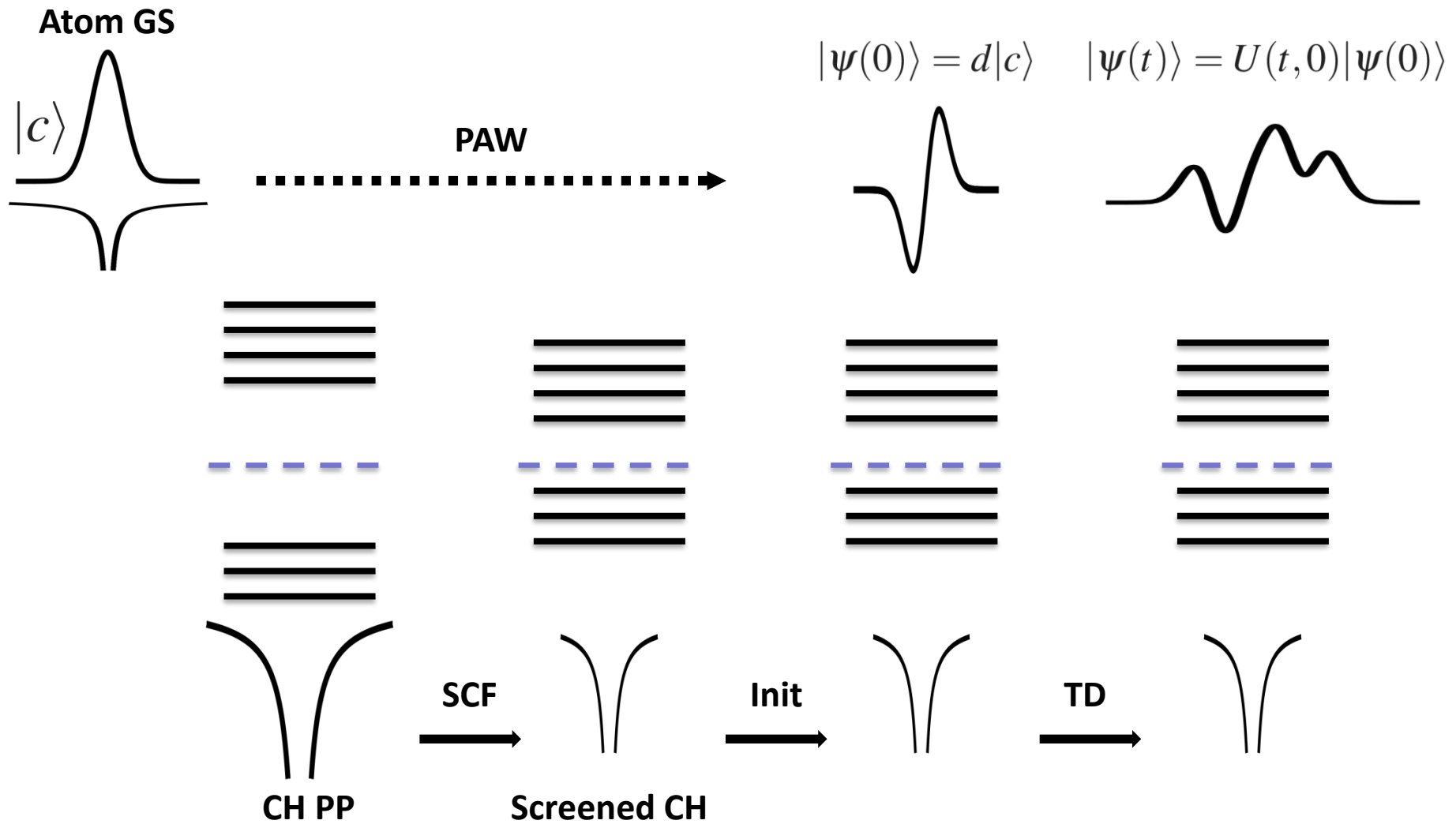
Crank-Nicolson

$$\bar{t} = t + \Delta/2 \quad \mathbf{U}(t + \Delta, t) = \frac{\mathbf{1} - \mathbf{S}^{-1} \mathbf{H}(\bar{t}) \Delta/2}{\mathbf{1} + \mathbf{S}^{-1} \mathbf{H}(\bar{t}) \Delta/2}$$

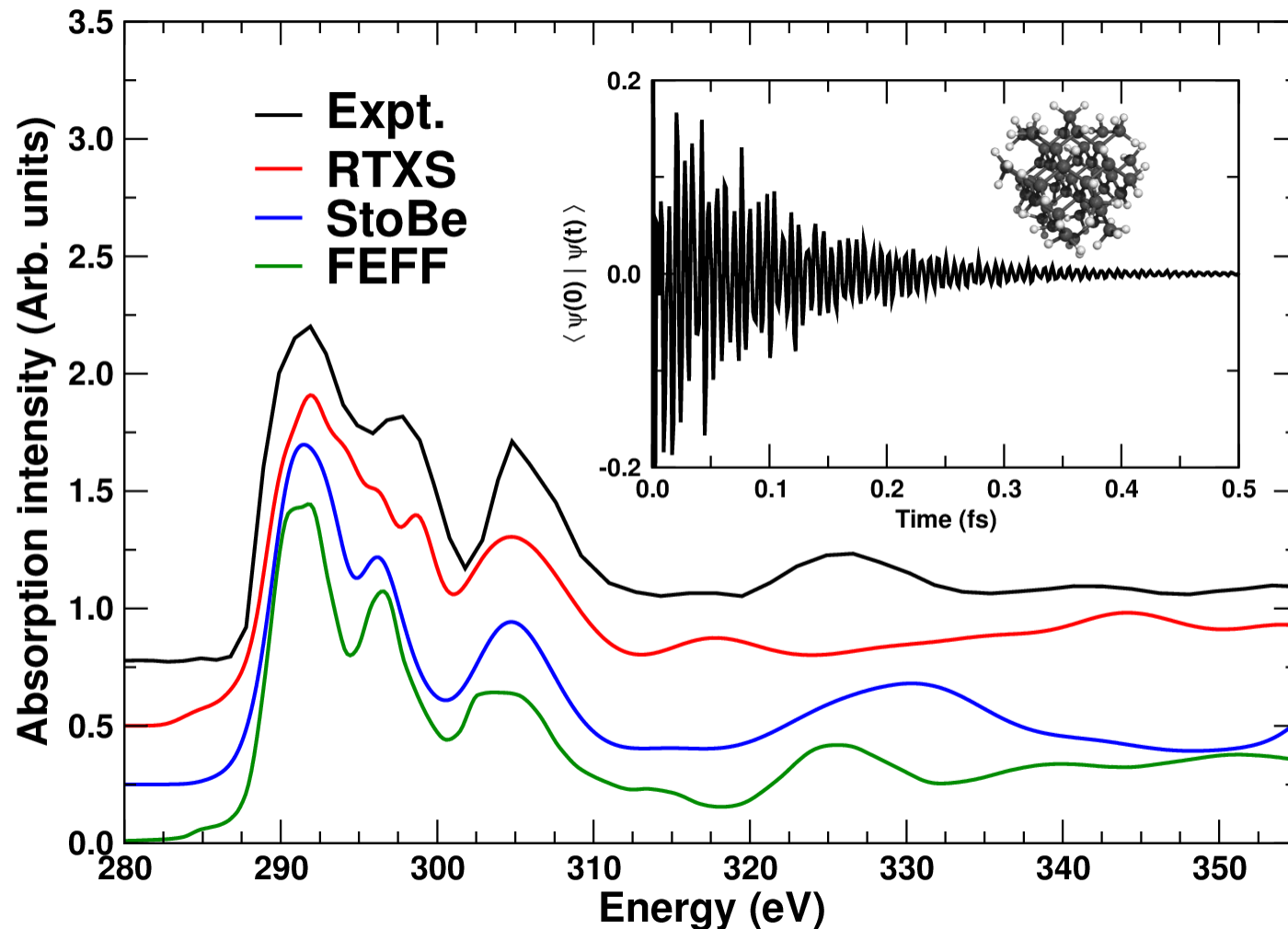
$$|\psi(t)\rangle = \sum_j |j\rangle c_j(t)$$

$$H_{jj'} = \langle j | h_H + v_{ch} + \Sigma | j' \rangle$$

RTXS in "Pictures"



C K-Edge XAS of Diamond ($C_{47}H_{60}$ cluster)



Expt: Fister et al., Phys. Rev. B 75, 174106 (2007)

Green's Functions and DFT/MD Approaches in X-Ray
Excited States and Structural Disorder

Real-Time Approaches for Optical and Core Response

Finite Temperature Green's Functions

Green's Functions at Finite-T

Motivation:

For **excited states & thermodynamics** at finite-T and extreme conditions (WDM, $T \sim T_F$) we need methods **beyond 0 K DFT**

Starting point: **Sum-rules** for energy and density

$$E(T)/N = \sum_k \int d\omega [\omega + \varepsilon_k] A_k(\omega) f(\omega) \quad n = \frac{1}{V} \sum_k \int d\omega A_k(\omega) f(\omega)$$

Key ingredient: **Spectral function** from Green's function

$$A_k(\omega) = (1/\pi) |\text{Im } G_k(\omega)|$$

Which Green's Function? GW vs Cumulant

GW

$$G(\omega) = G_0 + G_0 \Sigma G$$

$$\Sigma^{GW} = iGW \quad W = \epsilon^{-1} v$$

No vertex: $\Gamma = 1$

Cumulant*

$$G(t) = G_0(t) e^{C(t)}$$

$$C \sim / \text{Im } \Sigma^{GW} /$$

Implicit vertex

We choose the **cumulant** G : **Better** than GW

*Recent review: J. Zhou et al. J. Chem. Phys. 143, 184109 (2015)

Finite T Cumulant Green's Function

Similar to $T=0$ [1] with **implicit** temperature dependence [2]

$$G_k(t) = -i\theta(t)e^{-i\varepsilon_k^x t}e^{\tilde{C}_k(t)}$$

Retarded Green's function
Matsubara formalism

Natural separation into **independent particle** and **correlation** parts

$$\tilde{C}_k(t) = \int d\omega \frac{\gamma_k(\omega)}{\omega^2} (e^{-i\omega t} + i\omega t - 1) \quad \gamma_k(\omega) = \frac{1}{\pi} |\text{Im} \Sigma_k(\omega + \varepsilon_k)|.$$

FT Cumulant

[1] Kas et al., Phys. Rev. B 90, 085112 (2014)

[2] Kas et al., Phys. Rev. Lett. 109, 176403 (2017)

Thermodynamics from G

Thermodynamics of the homogeneous electron gas from the cumulant Green's function approach

T. Blanton, J. J. Kas, and J. J. Rehr
Dept. of Physics, Univ. of Washington Seattle, WA 98195
(Dated: January 5, 2018)

UW Preprint Jan 2018

Chemical potential $\mu(n, T)$: solution of $n(\mu, T) = n$

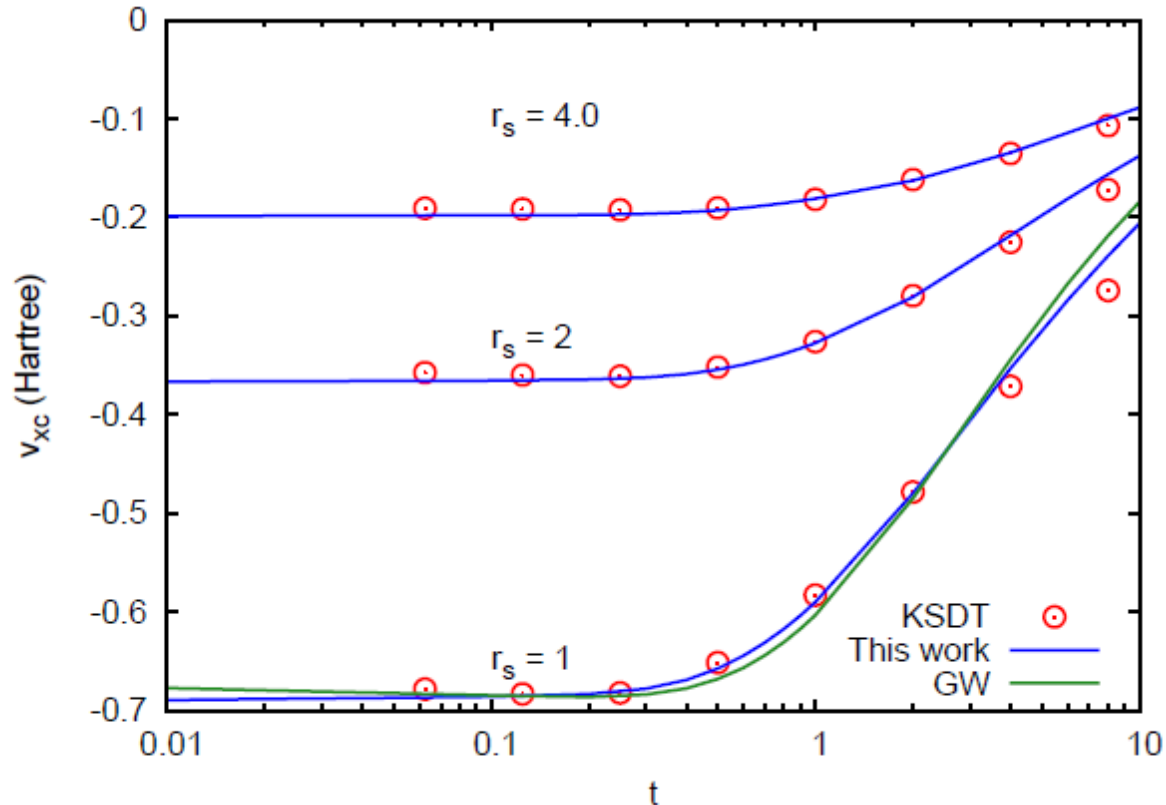
$$n = \frac{1}{V} \sum_k \int d\omega A_k(\omega) f(\omega) \quad f(\omega) = \frac{1}{e^{\beta(\omega - \mu)} + 1}$$

Can always separate into independent particle + xc parts:

$$\mu(n, T) = \mu_0(n, T) + \mu^{\text{xc}}(n, T)$$

Exchange-Correlation Potential

$$v_{xc}(T) \equiv \mu_{xc}(T) = \mu(T) - \mu_0(T)$$



Good agreement with finite T DFT functionals

Summary

Green's Functions:

Provide an **efficient** and **versatile** approach to:

X-Ray Excited States

Structural **Disorder**

Finite Temperature Effects

Real-Time Approaches:

RT-TDDFT:

Alternative to frequency-space simulations

Fast implementation, highly **adaptable**

Both **linear** and **non-linear** response

Access to **core** response

DFT/MD:

Crucial for systems with **dynamic** disorder/**bond breaking**

Real-time and finite temperature Green's function approaches for excited states, response functions, and thermodynamics

F. D. Vila, J. J. Kas and J. J. Rehr

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