

<sup>1</sup> **Towards Global Predictions of the Mode-1 Internal**  
<sup>2</sup> **Tide**

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5 A frequency and wavenumber analysis of satellite altimetry data for mode-  
6 1 internal tides over the central North Pacific ocean shows that these waves  
7 retain their spatial coherence over the 3500 km distance between the Hawai-  
8 ian Ridge and the Aleutians. The frequency-wavenumber tidal analysis re-  
9 solves signals consistent with the dispersion relation for internal tides. Waves  
10 propagating northward from Hawaii reach the Aleutians, and waves prop-  
11 agating southward from the Aleutians reach Hawaii. The result is consistent  
12 with the conclusion that mode-1 internal tides are likely predictable over most  
13 regions of the world's oceans. Accounting for the effect of stratification on  
14 the surface expression of the waves, mode amplitude, a quantity more di-  
15 rectly related to tidal energy, was derived. The internal-tide waves experi-  
16 ence little attenuation as they cross the North Pacific basin. The radiation  
17 of internal tides across ocean basins affects the spatial distribution of tidal  
18 energy available for ocean mixing.

## 1. Introduction

19 Recently, Dushaw et al. [2001] found that mode-1 internal tides were more predictable  
20 than previously appreciated. It appears likely that a significant aspect of oceanic variabil-  
21 ity, that associated with mode-1 internal tides, can be predicted over much of the world's  
22 oceans. This predictability should not be taken for granted. The amplitude and phase of  
23 internal tides have never been predicted at arbitrary points in space and time, although  
24 there have been hints of this property along altimeter tracks near topographic features  
25 such as the Hawaiian Ridge. There are numerous applications for predictions of internal  
26 tides, although, of course, these applications are scientific in nature and secondary to such  
27 applications as tidal predictions for navigation. We have only begun to exploit the ability  
28 to predict the mode-1 internal tides for oceanographic research.

29 Dushaw et al. [2011] used a combined frequency and wavenumber ( $\omega - \mathbf{k}$ ) harmonic  
30 analysis to derive maps of internal-tide harmonic constants from TOPEX/Poseidon (T/P)  
31 and Jason measurements of sea-surface height (SSH). These harmonic constants gave  
32 surprisingly robust and accurate predictions of mode-1 internal-tides near the Hawaiian  
33 Ridge, to at least  $40^\circ\text{N}$  [Dushaw et al. 1995; Rainville et al. 2009], and in the western  
34 North Atlantic [Dushaw 2006; Dushaw et al. 2011]. Predictions are limited to the first  
35 internal tide mode, since that mode has the dominant expression in sea surface height,  
36 a long wavelength ( $\sim 150$  km) and a fast phase speed ( $\sim 3.5$  m/s). The predictability of  
37 higher modes is unknown. The amplitude and phase of predictions derived from altimetry  
38 closely matched the time series of historical in situ measurements by acoustic tomography,

39 even though some of the historical data was obtained a decade or more before the altimetry  
40 data used to make the prediction.

41 Developing accurate predictive models for the internal tide will likely be a long process,  
42 however. The comparisons described in detail by Dushaw et al. [2011] demonstrate that  
43 these predictions are possible. It is now a matter of (1) developing procedures to hone the  
44 robustness and accuracy of the predictions, and (2) determining the limitations of those  
45 predictions. The regions where internal tide estimates can be tested by tomographic  
46 data, the North Pacific and western North Atlantic, are fairly benign, with fairly weak  
47 surface currents and mesoscale variability. It may be, or indeed seems likely, that the  
48 predictability apparent in these regions will not be so good in other regions. From the  
49 global perspective, the task at hand is to assess the internal-tide predictability, try to  
50 adapt the model employed in the tidal analysis of SSH for the specific conditions of each  
51 region to best extract predictable internal tides, or to understand why the predictability  
52 fails.

53 Dushaw et al. [2011] focussed on the tidal analysis over regional areas to derive pre-  
54 dictions in specific locations for comparing to the available in situ data. In this paper,  
55 the analysis is applied to the central North Pacific between the Hawaiian Ridge and the  
56 Aleutian Islands. The resulting maps show that the internal tide waves retain their co-  
57 herence as they cross ocean basins. For these waves, the time scale associated with this  
58 3500-km distance is about 12 days. In addition, after properly accounting for the effects  
59 of varying stratification and latitude on the sea-surface-height expression of the waves,  
60 the dissipation of the waves is shown to be remarkably weak.

## 2. Internal tide propagation between Hawaii and the Aleutians

61 As discussed by Dushaw et al. [2011], applying an  $\omega - \mathbf{k}$  analysis to satellite data  
62 over a relatively large area is problematic because it is difficult to formulate an adequate  
63 wavenumber model for large areas. Indeed, it is likely the wavenumber content is not  
64 stationary across large areas. To work around this issue, the large area between Hawaii  
65 and the Aleutians was divided into subdomains of  $11 \times 11^\circ$ , each overlapping adjacent  
66 subdomains by  $2^\circ$ . The  $\omega - \mathbf{k}$  analysis was applied to each subdomain independently.  
67 The solution over the entire domain was then constructed by merging the results from  
68 each subdomain using a simple cosine taper over the  $2^\circ$  overlap (Fig. 1a). Ten years of  
69 T/P+Jason altimeter data were used for the analysis, together with tandem track data  
70 [Dibarboure et al. 2009]. To estimate just the waves emanating northward from the  
71 Hawaiian ridge, the tidal model for Figure 1a assumed waves propagating only in the gen-  
72 erally northward direction. The tidal fit therefore excludes southward propagating waves.  
73 The  $\omega - \mathbf{k}$  analysis obtains a solution consistent with only the assumed wavenumbers,  
74 and excludes variability not consistent with the internal tide dispersion relation, such as  
75 the mesoscale. For the  $M_2$  tidal constituent, coherent waves radiating from Hawaii to  
76 the Aleutians and the Gulf of Alaska are apparent, giving a stark demonstration of the  
77 advantage of the  $\omega - \mathbf{k}$  analysis. Zhao and Alford (2009) computed energy fluxes for north-  
78 ward and southward waves that showed internal-tide energy radiating across the North  
79 Pacific basin. Here, these waves have been resolved to the extent that it is apparent that  
80 internal-tide waves cross ocean basins without much attenuation or loss of coherence.

81 A similar tidal analysis using only southward propagating waves found internal tide  
82 waves that propagate from the Aleutian Islands to the Hawaiian Ridge (Fig. 1b). Waves  
83 generated at the Ridge and radiating southward are also apparent. An optimal tidal anal-  
84 ysis includes wavenumbers in both directions (results not shown). The addition of both  
85 northward and southward propagating waves produces an internal-tide field characterized  
86 by ubiquitous standing waves [Zhao and Alford 2009], [Zhao et al. 2009], [Dushaw et  
87 al. 2011]. The existence of these standing waves in the central North Pacific, formed by  
88 the interference of waves from Hawaii and waves from Alaska, is a telling indicator of the  
89 temporal coherence of these waves, inasmuch as an interference pattern is particularly  
90 sensitive to the phase variations of the underlying waves.

### 3. Sea-surface height and stratification

91 Dushaw et al. [2011] showed that stratification has a large influence on the surface  
92 expression of the internal tide. The dependence of surface expression on stratification can  
93 be illustrated by calculating the mode-1 displacement modes using the 2009 World Ocean  
94 Atlas [Antonov et al. 2010; Locarnini et al. 2010] and assuming a constant internal dis-  
95 placement over the world's oceans (Fig. 2). The calculation assumed a flat-bottom ocean  
96 of 5500 m depth to determine the effects of stratification alone. A proper calculation of  
97 the relation between mode amplitude and SSH would be a solution for the modes using the  
98 local depth and stratification, followed by a normalization to ensure equal energy density  
99 of the modes for a unit mode-1 amplitude. The effects of varying depth and energy are  
100 small over most of the oceans, however, since the ocean basins have fairly constant depths  
101 of 5000-6000 m. In any case, it is apparent that for a given internal amplitude, the surface

102 expression of internal tides in tropical latitudes is generally much greater than in higher  
103 latitudes. This effect has led to general conclusions concerning the decay properties of the  
104 internal (e.g., *Tian et al. [2008]*), but accounting for the effects of stratification changes  
105 notions of the decay of the internal tide considerably. Calculating internal displacement,  
106 more directly related to internal-tide energy density than SSH, shows that the mode-1 in-  
107 ternal tide has a much larger decay scale than previously appreciated (Fig. 3). Indeed, for  
108 cases such as waves radiating northward from the Hawaiian Ridge, cylindrical spreading  
109 must also be accounted for when estimating a decay length scale.

#### 4. Comments on the calculation of internal-tide energy flux

110 Although it is a common procedure (e.g., Ray and Cartwright [2001], *Dushaw [2002]*,  
111 *Zhao and Alford [2009]*), calculating energy flux at single points of a field resulting from  
112 multiple interfering waves, or by fitting monochromatic plane waves, is problematic, if for  
113 no other reason than it provides an erroneous physical picture. For example, one of the  
114 reasons the Hawaiian Ocean Mixing Experiment (HOME) [*Pinkel et al. 2000; Carter et*  
115 *al. 2008; Rainville et al. 2009*] focussed on Kaena Ridge was because of the maps showing  
116 “beams” of internal-tide energy apparently emanating from there. It is now clear that  
117 much of the energy of those “beams” did not emanate from the Kaena Ridge. Rather, the  
118 “beams” are formed by the superposition of multiple waves emanating from a wide stretch  
119 of the Hawaiian Ridge [*Rainville et al. 2009, Zhao and Alford 2009*]. Also, the energy flux  
120 calculated based on the interference “beams” suggests a decay length scale that bears little  
121 resemblance to the actual decay scale of mode-1 internal tides. The energy flux derived  
122 from “beams” diminishes when the waves are no longer in constructive interference (c.f.,

123 Fig. 1a). Such “beams” do not have a real physical meaning, other than that the tidal  
124 amplitude there is large because of constructive interference (Fig. 4).

125 Two rigorous, equivalent approaches to estimating the mode-1 energy flux of compli-  
126 cated wave fields are proposed. The energy flux is fundamentally defined as  $(\langle p'u' \rangle, \langle p'v' \rangle)$ .  
127 In the context of the interference patterns of the internal tide around Hawaii, some care  
128 must be given to defining what is meant by the brackets  $\langle \cdot \rangle$ . This averaging implies an  
129 average over time to remove the oscillations of energy flux with time, an average over  
130 depth to remove the oscillations of energy flux with depth, and, by analogy, an average  
131 over area to remove the oscillations of energy flux with area. This latter average is most  
132 often omitted, giving the common figures of beams of large energy apparently emanating  
133 from such locations as the Kaena Ridge. Without the area average, expressions such as  
134  $\langle p'u' \rangle$ , with  $p'$  and  $u'$  expanded as grand summations over frequency, mode, and wavenum-  
135 ber, retain cross terms of wavenumbers which have no clear physical meaning. With such  
136 averaging, the energy flux reduces to a simple summation of the square of the individual  
137 components, as required.

138 As an alternative to the averaging implied by the brackets, the variability can be sep-  
139 arated in frequency, mode, and wavenumber (time, depth, and area) manually. A tidal  
140 analysis can separate the signals of the constituent frequencies, a modal decomposition  
141 can separate the signals of individual modes, and a wavenumber decomposition can sep-  
142 arate the signals of individual wavenumbers. With this separation, no averaging  $\langle \cdot \rangle$  is  
143 required, since the individual components of the tidal variability can be directly added.  
144 The frequency and wavenumber decompositions are a direct product of the  $\omega - \mathbf{k}$  tidal

145 analysis. The tidal signals derived from SSH discussed here are dominated by mode 1  
146 and consistent with the mode-1 dispersion relation, giving a natural isolation of the first  
147 mode.

## 5. Discussion

148 Dushaw et al. [2011] concluded that the mode-1 internal tide is likely predictable, in  
149 the traditional tidal sense, over most regions of the world's oceans. Temporally incoherent  
150 contributions to mode-1 internal tides appeared to be minimal, at least in the North Pacific  
151 and western North Atlantic oceans, and the waves showed very little attenuation. Here,  
152 consistent with those conclusions, the mode-1 internal tides are found to have a spatial  
153 coherence that extends across the North Pacific basin. Further, after correcting for the  
154 effect of local stratification on the surface expression of the internal tides and allowing for  
155 cylindrical spreading, little decay of the mode-1 amplitude is evident over the 3500 km  
156 distance between the Hawaiian Ridge and the Aleutians.

157 One aim of this project is to provide accurate predictions of the internal-tide signatures  
158 of sea-surface height, so that this noise element may be removed from other observations  
159 such as satellite observations by the Surface Water and Ocean Topography (SWOT) mis-  
160 sion. Such predictions also will require an assessment of their reliability, depending on the  
161 ocean region of interest. It seems unlikely that internal tides in all areas of the world's  
162 oceans will be as predictable as they seem to be in the North Pacific. A quantitative as-  
163 sessment of this predictability over the world's oceans is required. High-resolution models  
164 or state estimates such as are available from the Estimating the Circulation and Climate  
165 of the Ocean (ECCO2) program [Menemenlis et al. 2008] may give an accurate estimate

166 for the environment through which internal tides propagate, hence may help identify the  
167 processes that disrupt internal-tide coherence (see also Arbic et al. [2010]). Including all  
168 tidal constituents, the signals of internal tides in SSH have typical amplitudes of 3 – 5 cm  
169 in most regions of the ocean.

170 Unlike the basic premise of HOME, the tidal energy budget and its relation to deep-  
171 ocean mixing is not a local problem. Internal tide waves that carry most of the energy lost  
172 from the barotropic tide [Carter et al. 2008] at any location can cross ocean basins with  
173 little dissipation to deposit their energy on distant coastlines. Meanwhile, even at remote  
174 locations such as the Hawaiian Ridge, radiated internal-tide energy can arrive from distant  
175 regions. When it arrives at coastlines or shoaling topography, internal-tide energy either  
176 dissipates or reflects. These energy considerations determine the spatial distribution of  
177 tidally-driven deep ocean mixing, an essential aspect of general circulation models.

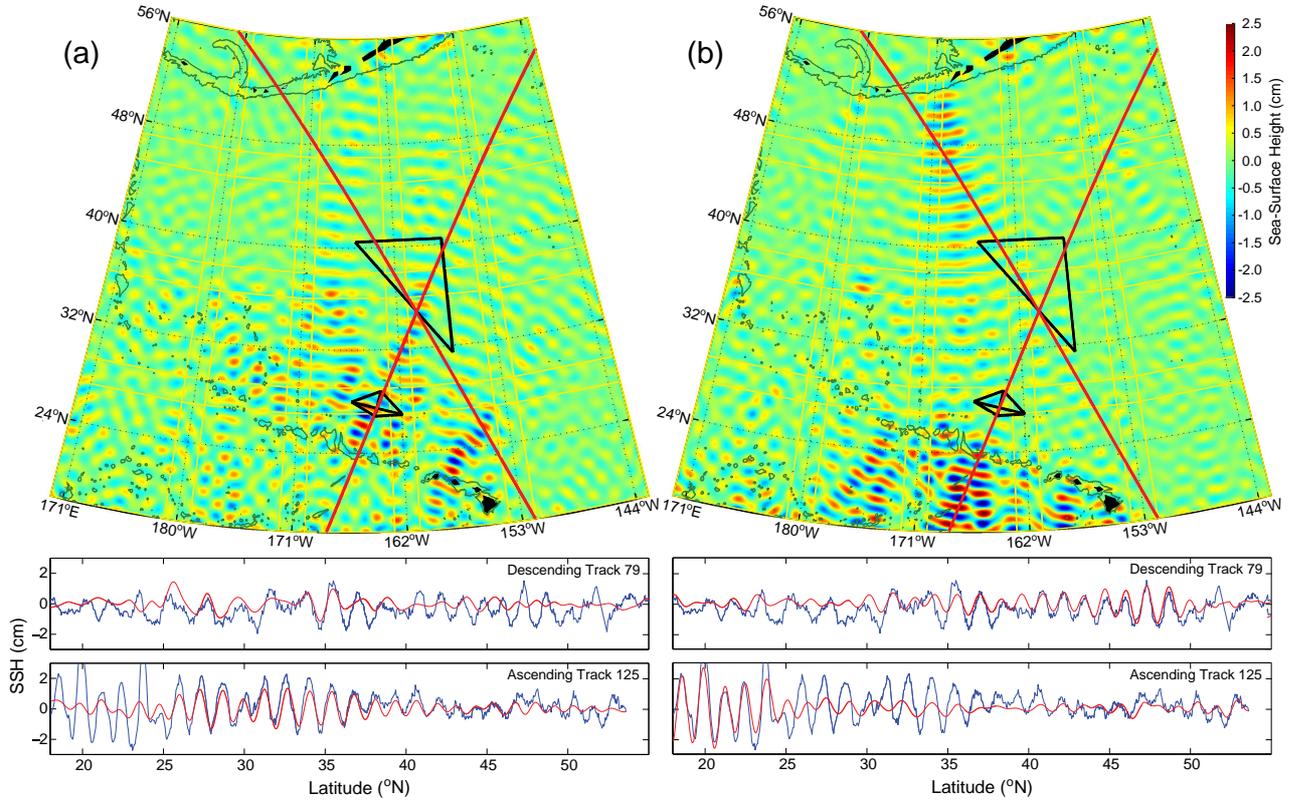
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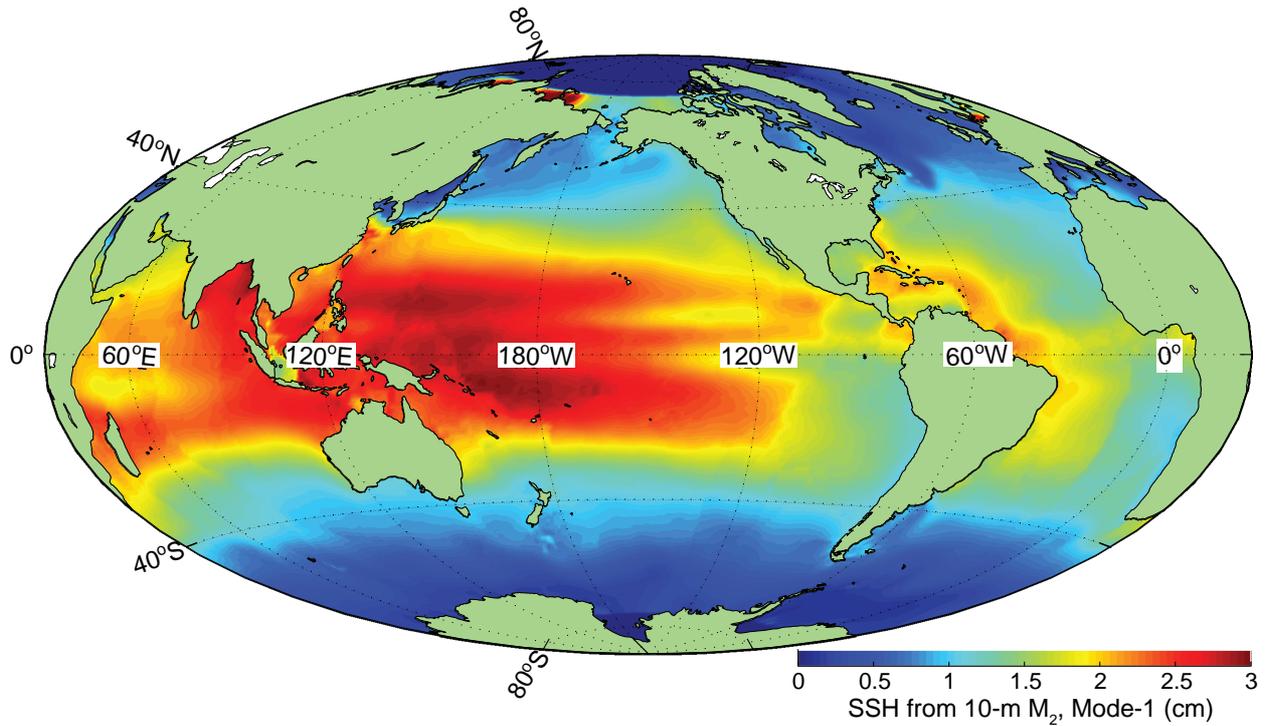
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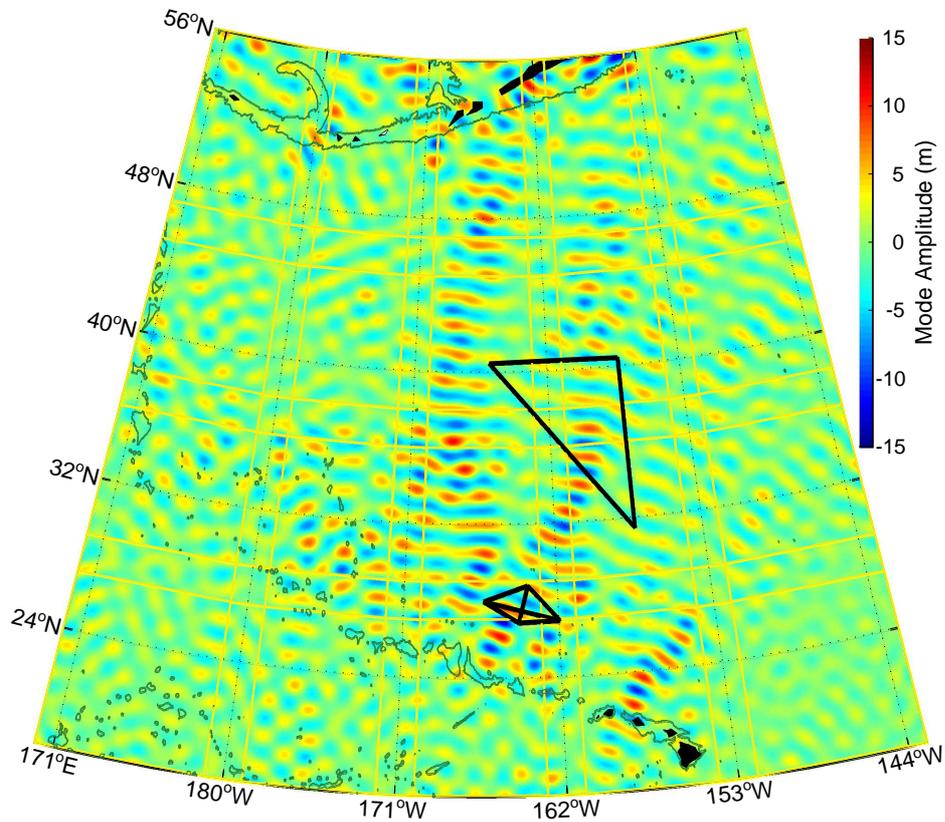
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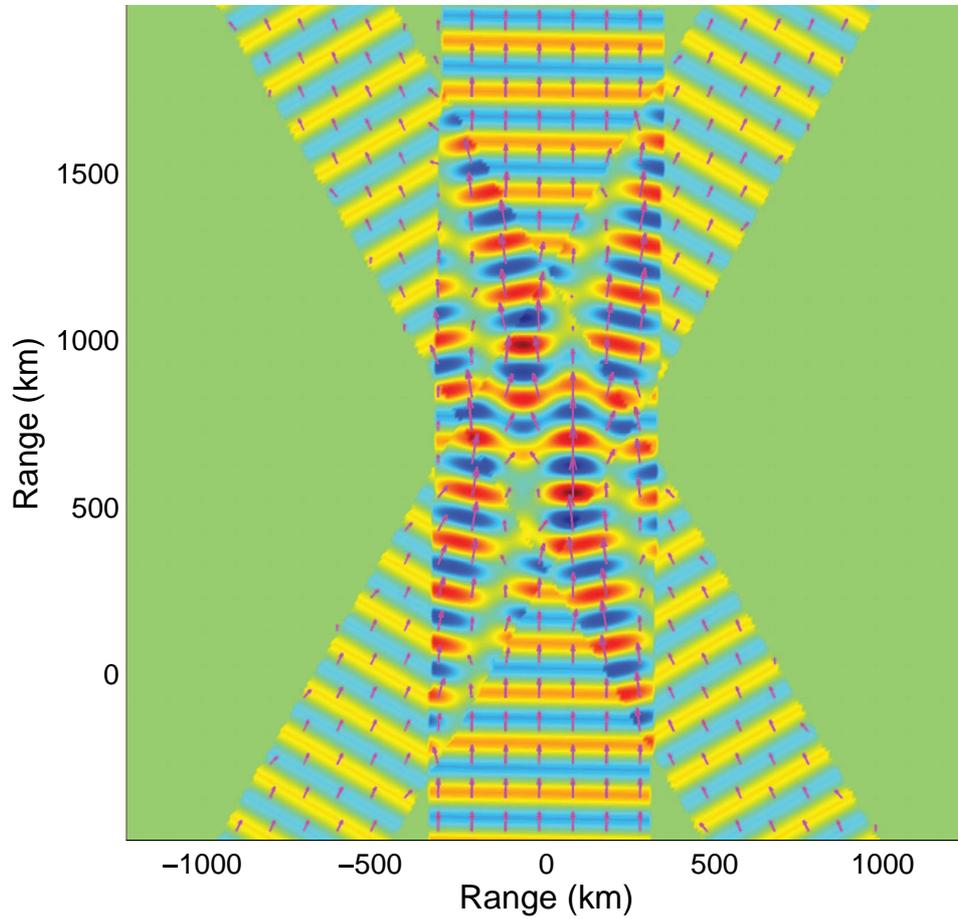
**Figure 1.** (a) The cosine component of the harmonic constants ( $A \cos(G)$ ) for the  $M_2$ , mode-1 internal tide derived from a  $\omega - \mathbf{k}$  analysis of T/P+Jason SSH assuming wavenumbers in the northward direction. The discrete yellow lines indicate the boundaries of the independent subdomains, overlapping by  $2^\circ$ . The RTE87 triangle [Dushaw *et al.* 1995] and HOME diamond tomography arrays [Rainville *et al.* 2009] are indicated. (b) Assuming wavenumbers of general southward direction. Lambert azimuthal equal area projection. (Bottom panels) Comparison of point-wise (blue) and  $\omega - \mathbf{k}$  (red) harmonic constants along tracks 79 and 125, indicated by the red lines on the top maps.



**Figure 2.** Amplitude of SSH in cm of a mode-1 internal tide at  $M_2$  frequency with 10-m internal displacement amplitude over the world's oceans derived from the 2009 World Ocean Atlas. The surface expression of the internal tide depends on the local stratification, hence local conditions must be accounted for when calculating energy or mode amplitude from a measurement of SSH.



**Figure 3.** The mode amplitude derived from the map of SSH assuming northward wavenumbers (Fig. 1a). Internal tides experience quite weak decay as they propagate across ocean basins. With the mode function normalization employed here, “Mode Amplitude” is roughly equal to the maximum internal displacement of the mode in meters.



**Figure 4.** Simple illustration of the problem of deriving point-wise energy flux vectors from a field comprised of interfering waves. Where these waves form a complicated interference pattern at their intersection, a complicated field of energy flux vectors is obtained. In this case, there are only three energy flux vectors, one for each monochromatic wave.